



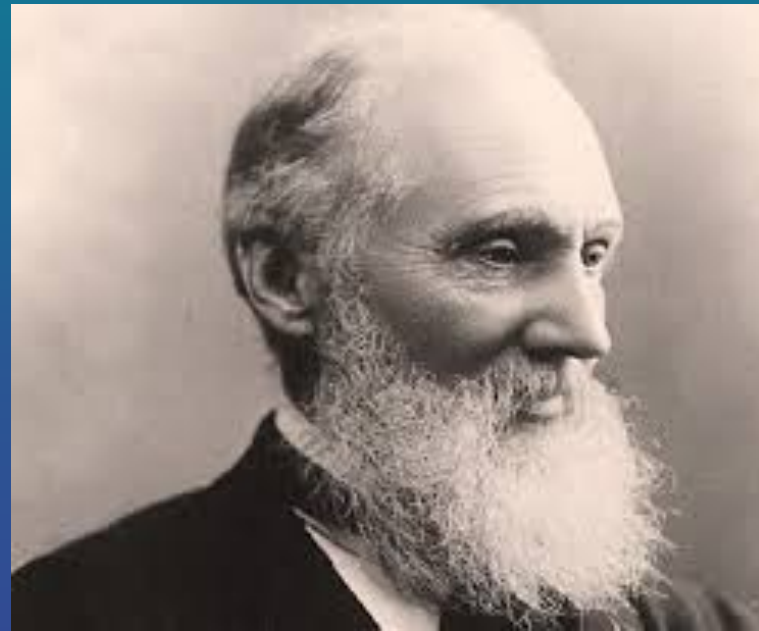
Sharif Quantum Information Group

Topology and Quantum Computation

Vahid Karimipour,
Sharif University of Technology, Iran.

Sharif Colloquium
1395

The origins of Knot Theory



Kelvin

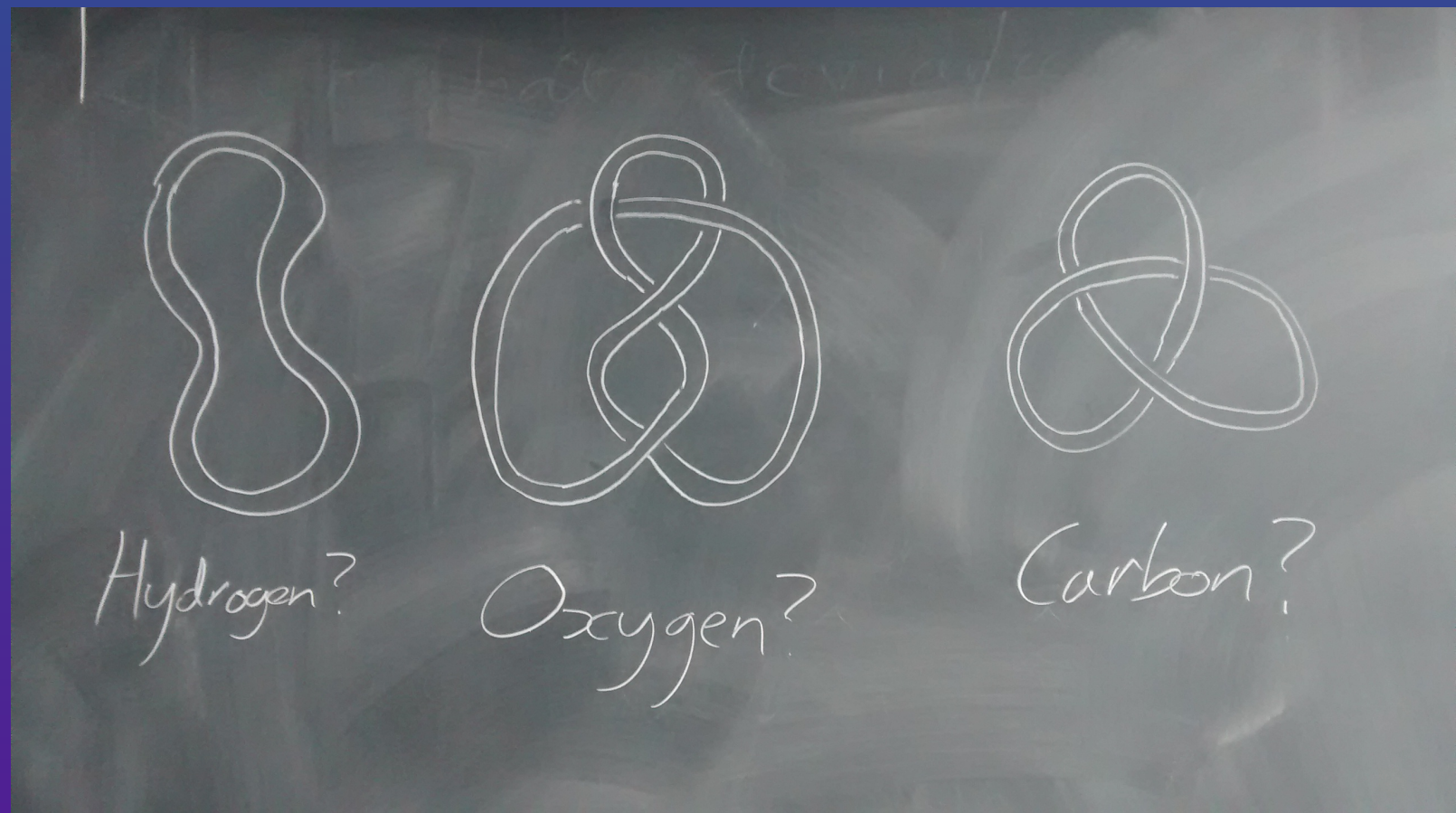
Vortex lines in fluids are stable!

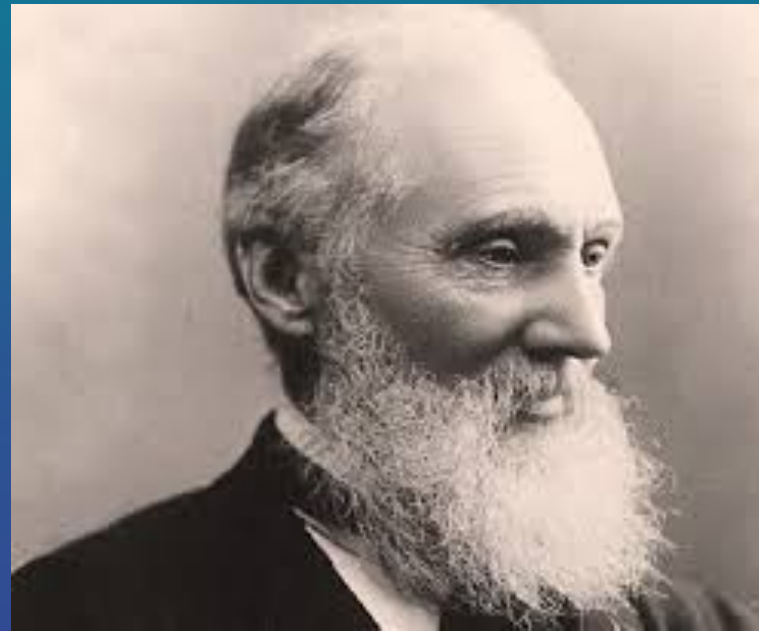
Atoms are Knots in Ether!

They are stable and take many forms.



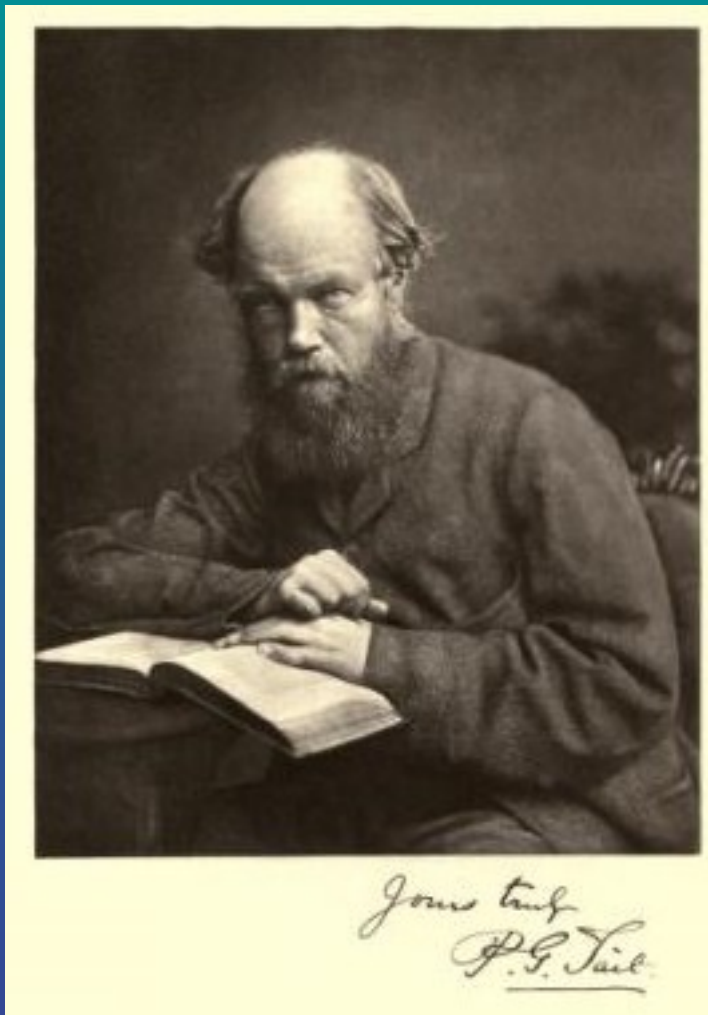
Maxwell





Kelvin

Not interesting at all.



Peter Tait

Very interesting.



Maxwell

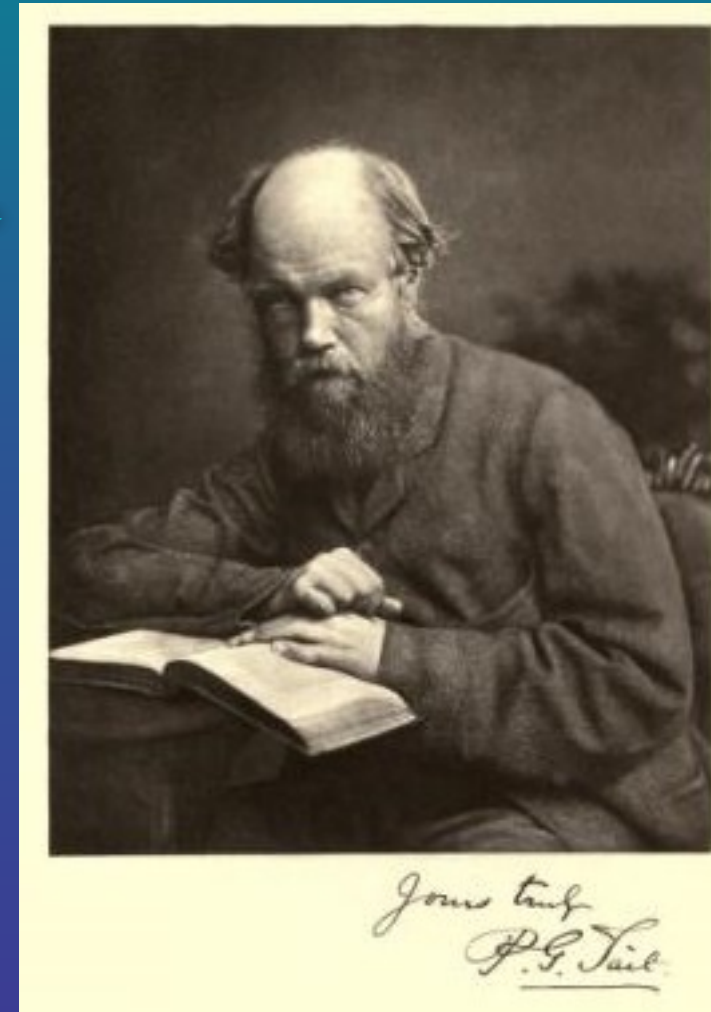


After 10 Years!



Kelvin

Very interesting.



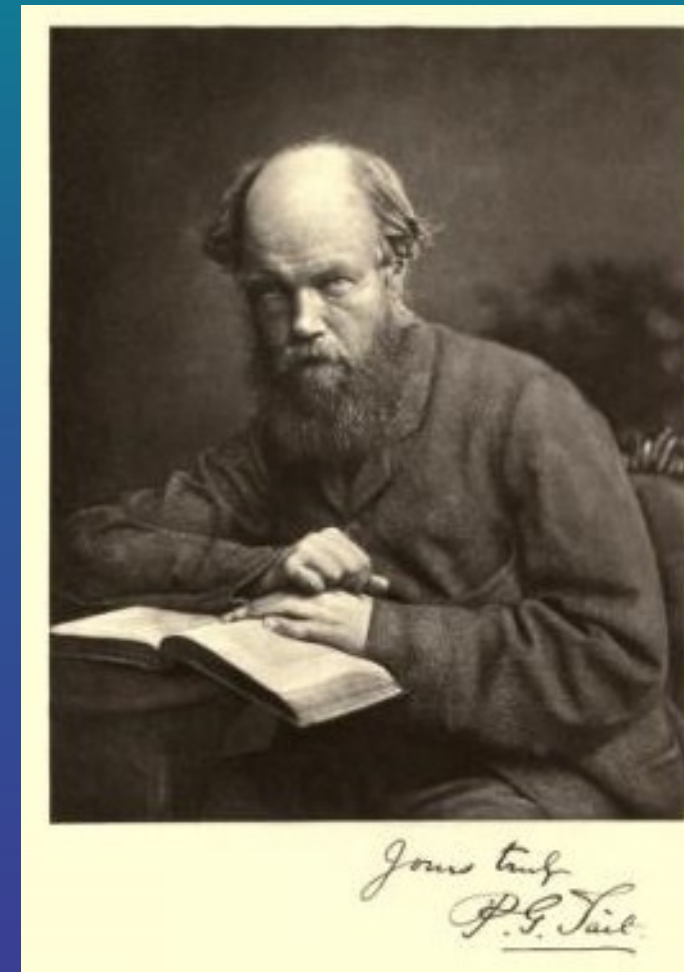
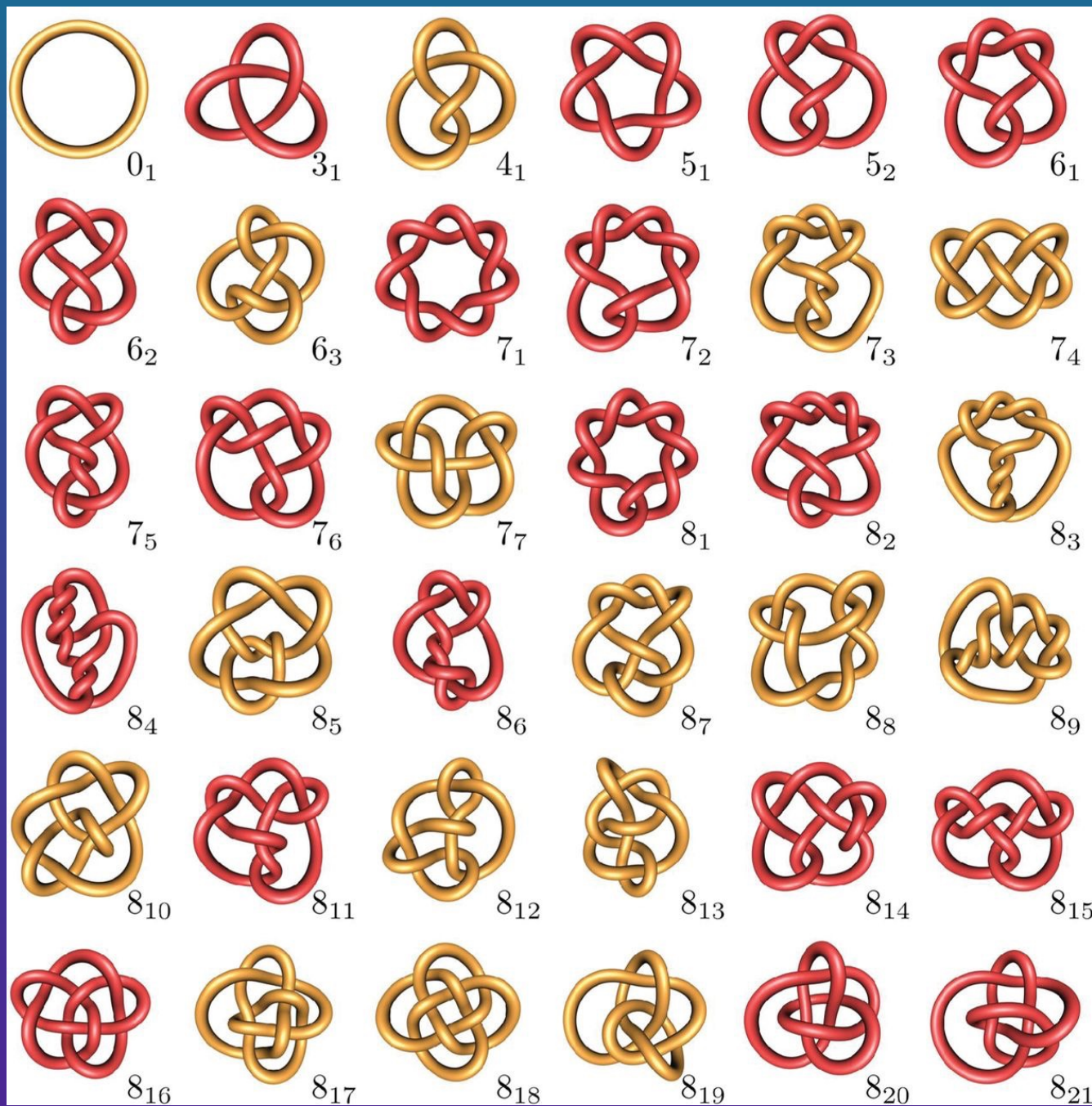
Not interesting at all.

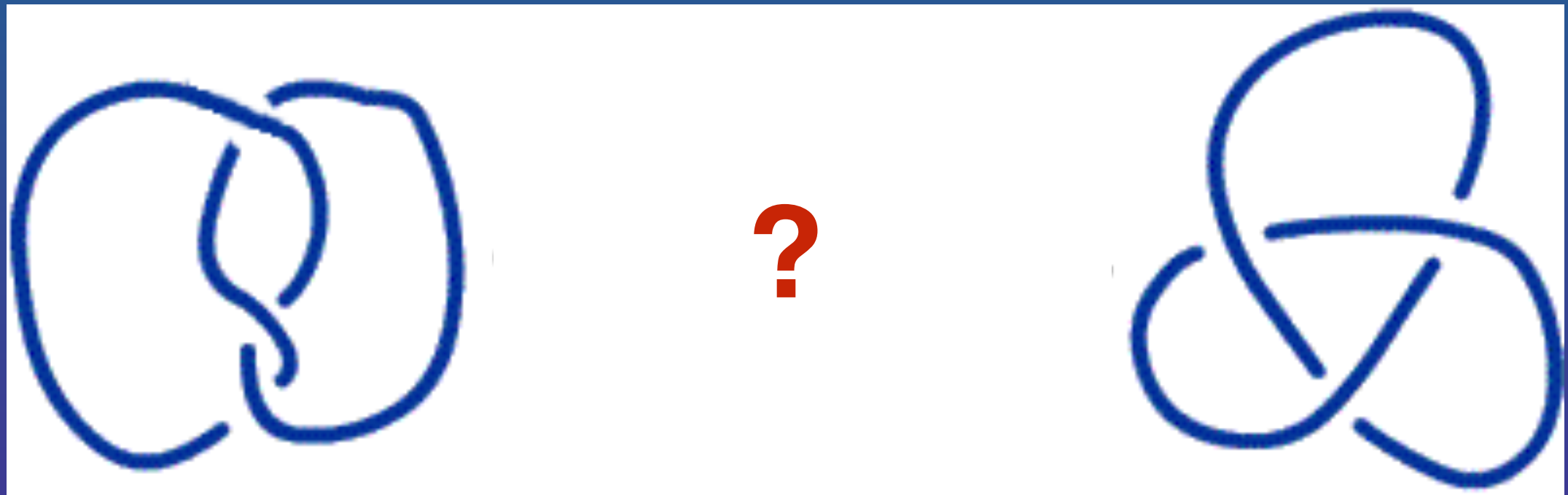


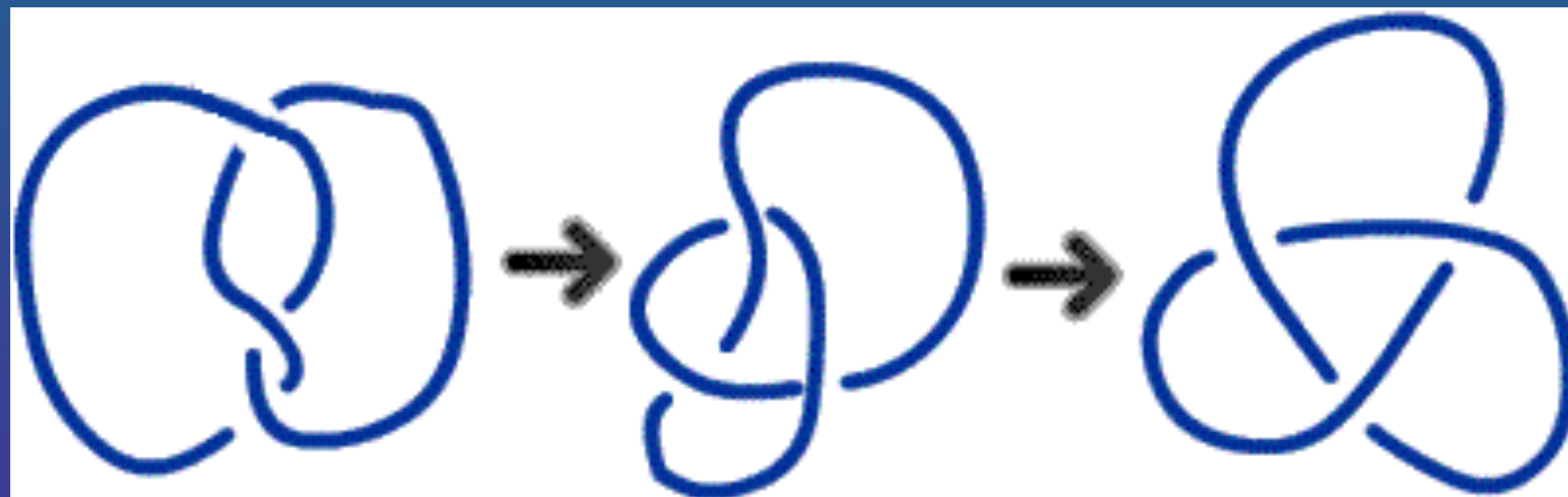
Maxwell



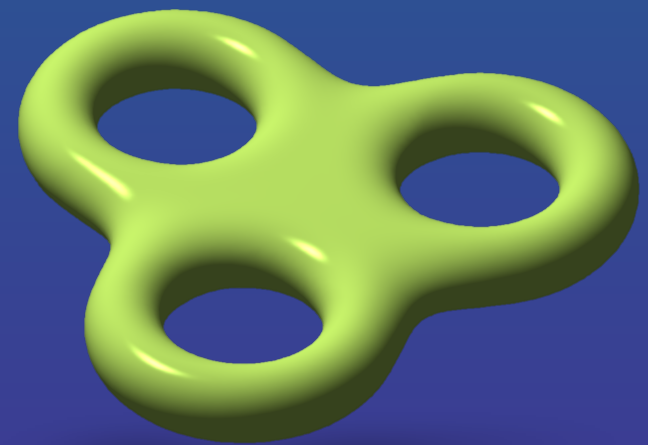
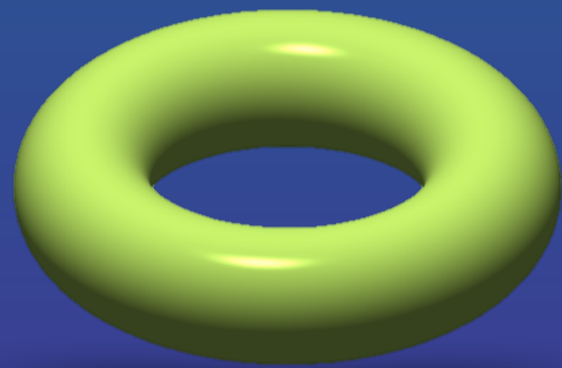
Knots and periodic table of elements



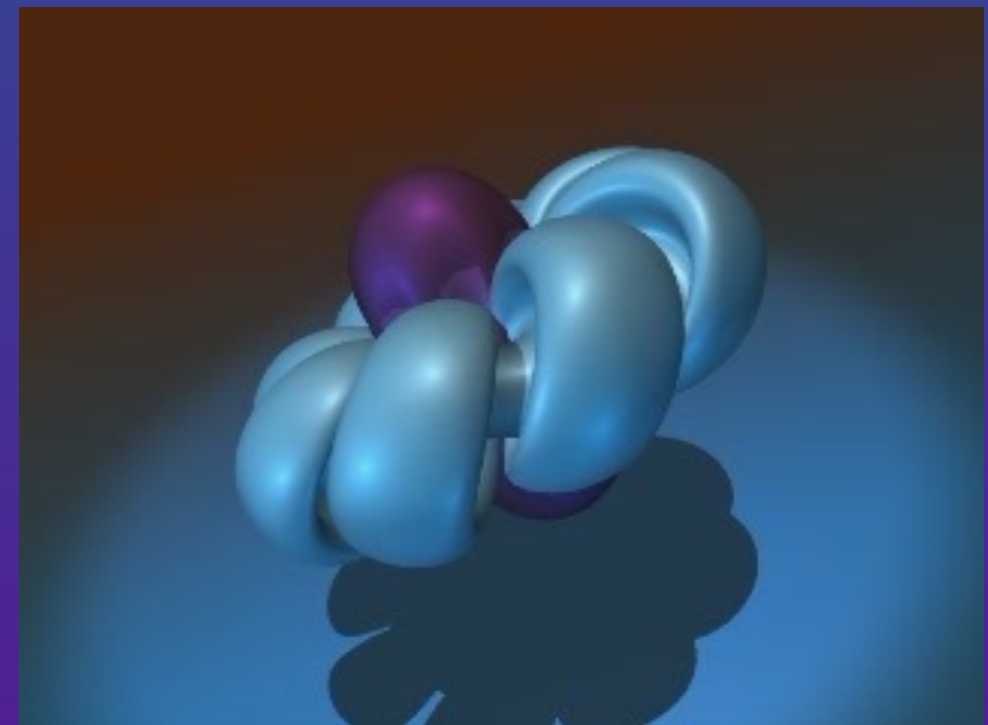
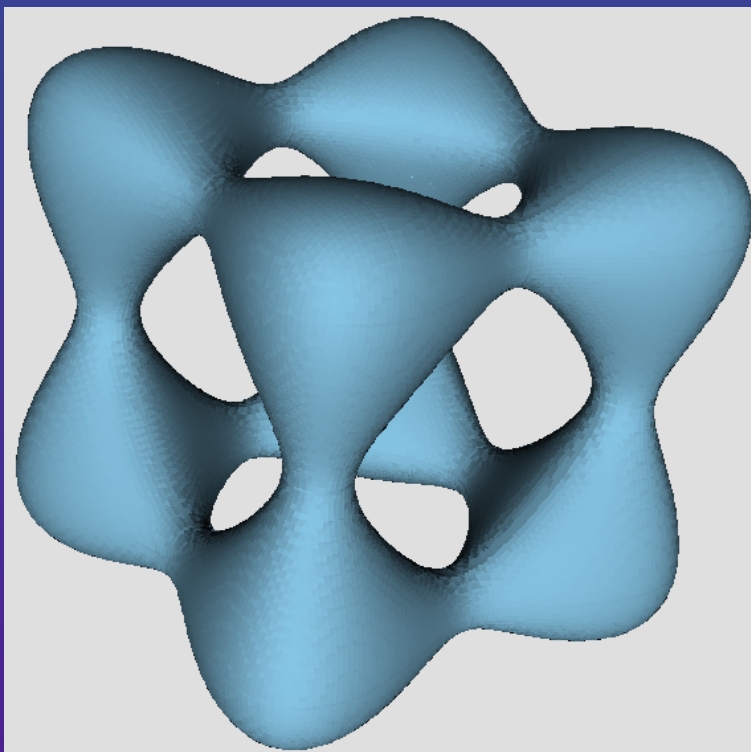
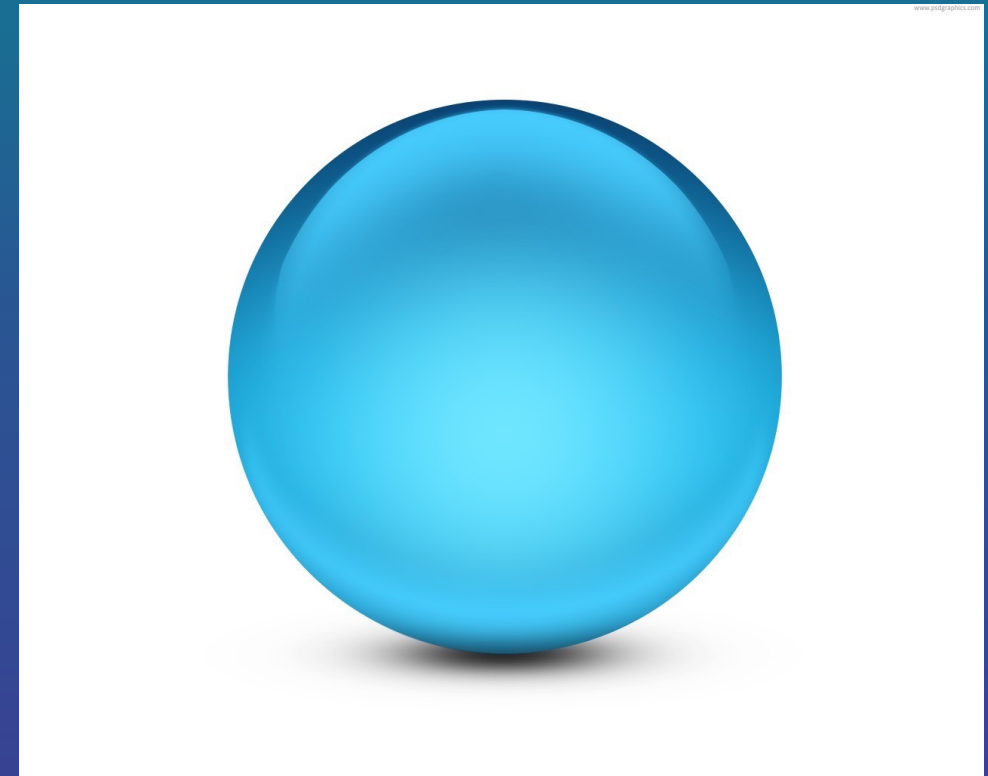
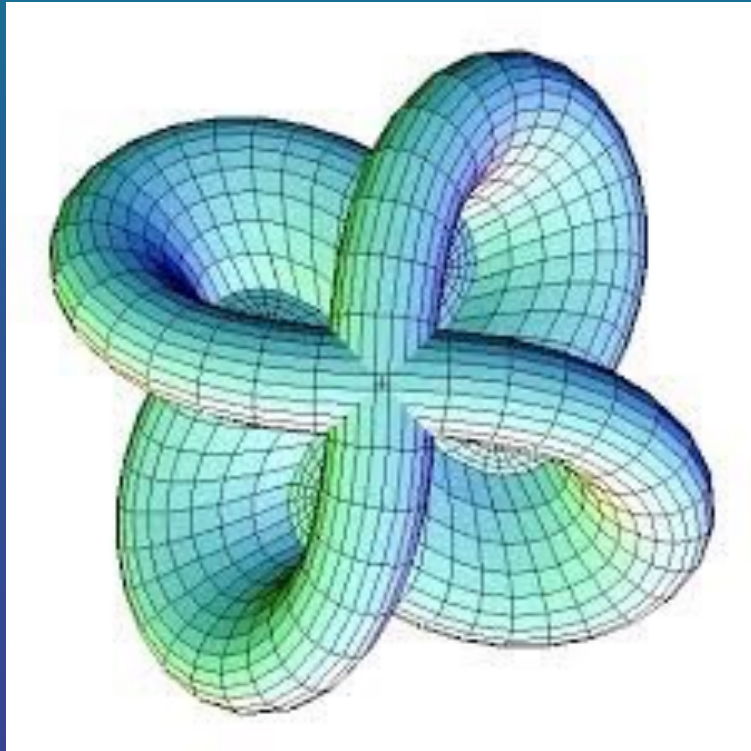




Two Dimensional Manifolds

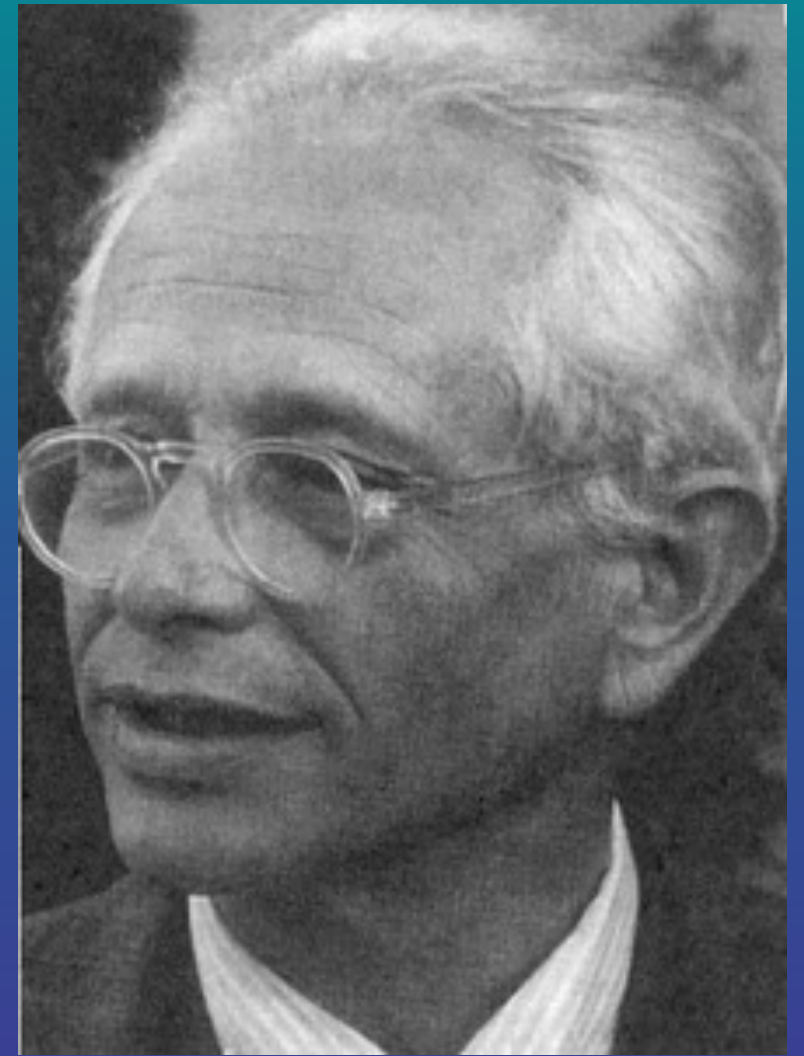


Three Dimensional Manifolds?



Dehn Surgery

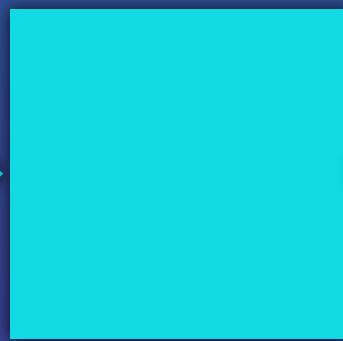
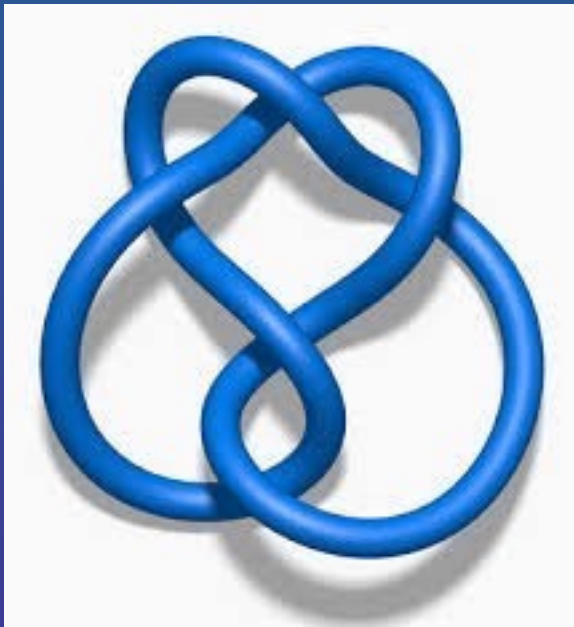
$$S_3 \longrightarrow M$$



Max Dehn

Knot Invariants

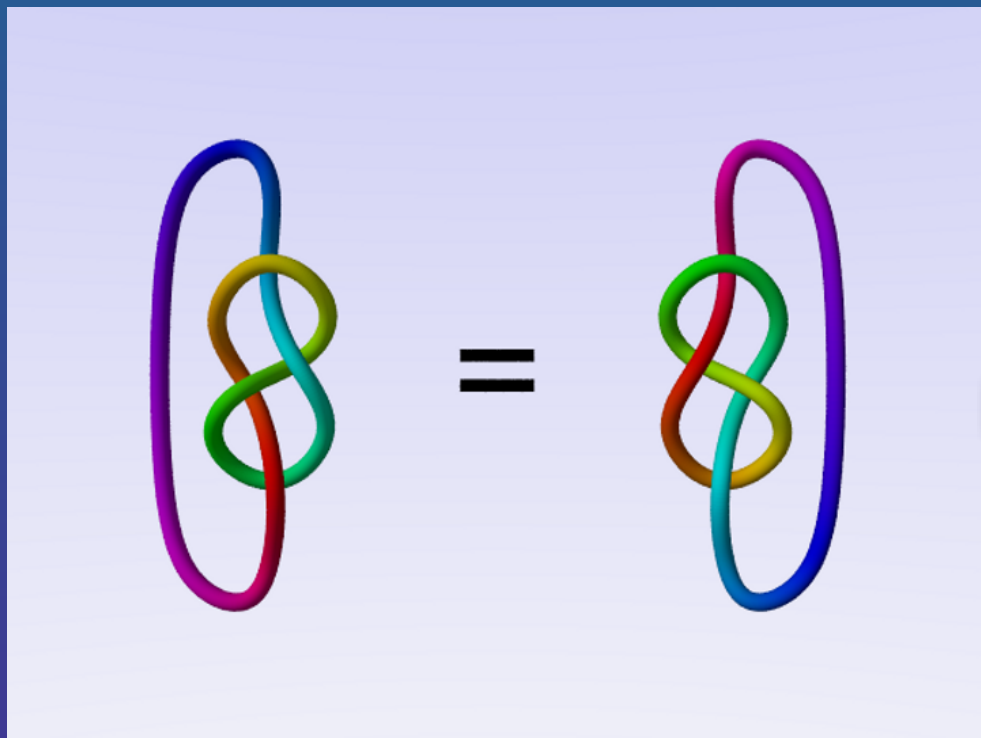
K



$W(K)$

Knot Invariants

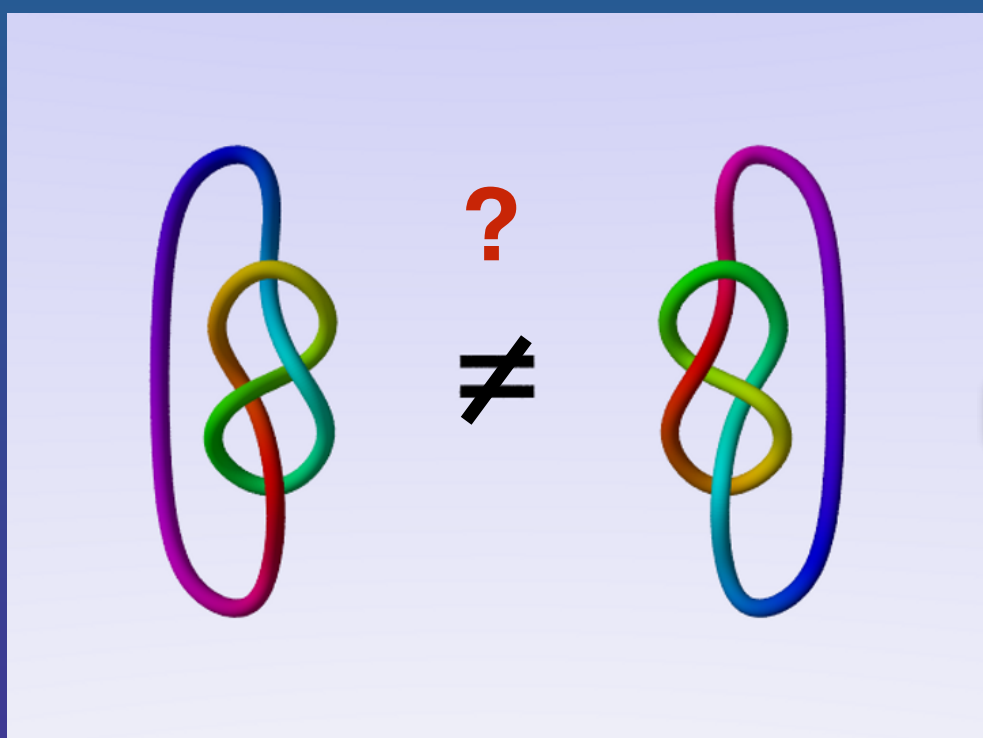
$$K \sim K'$$



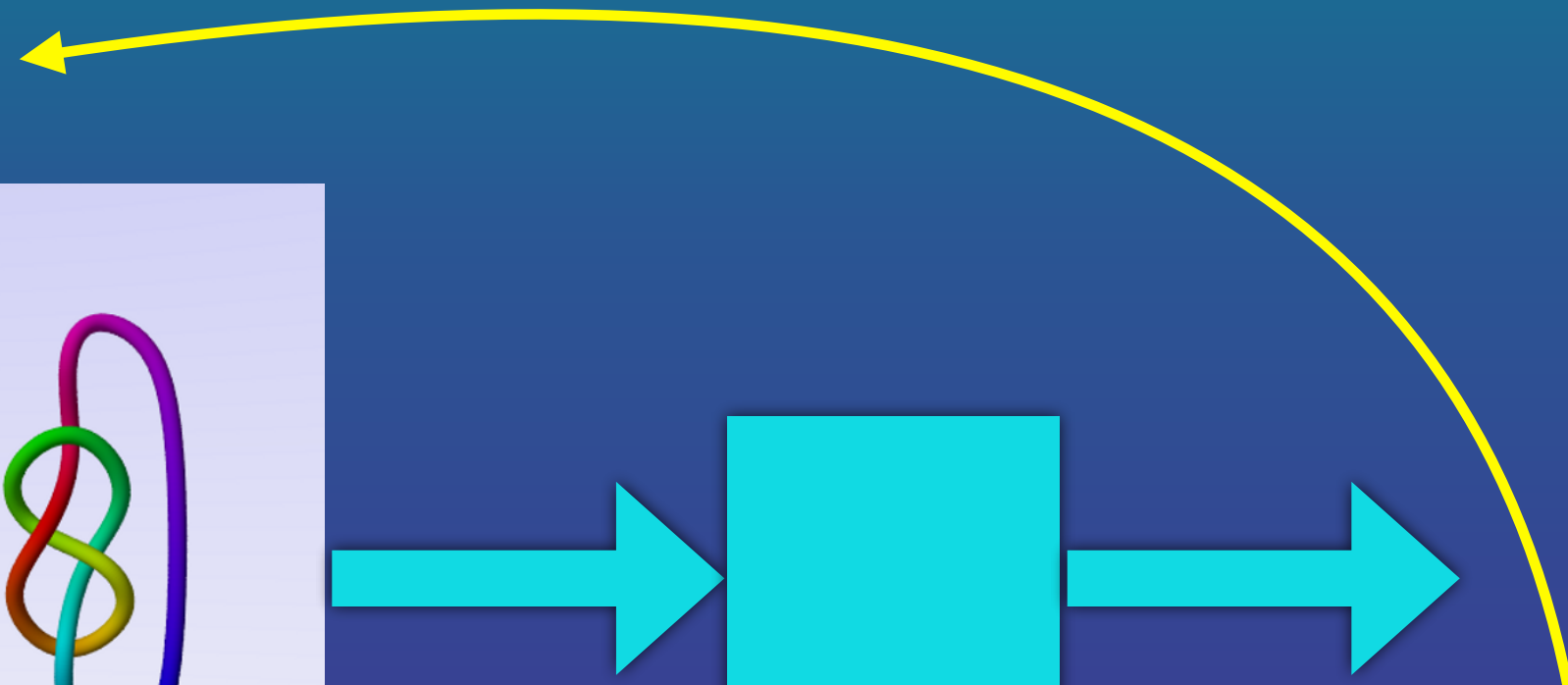
$$W(K) = W(K')$$

Knot Invariants

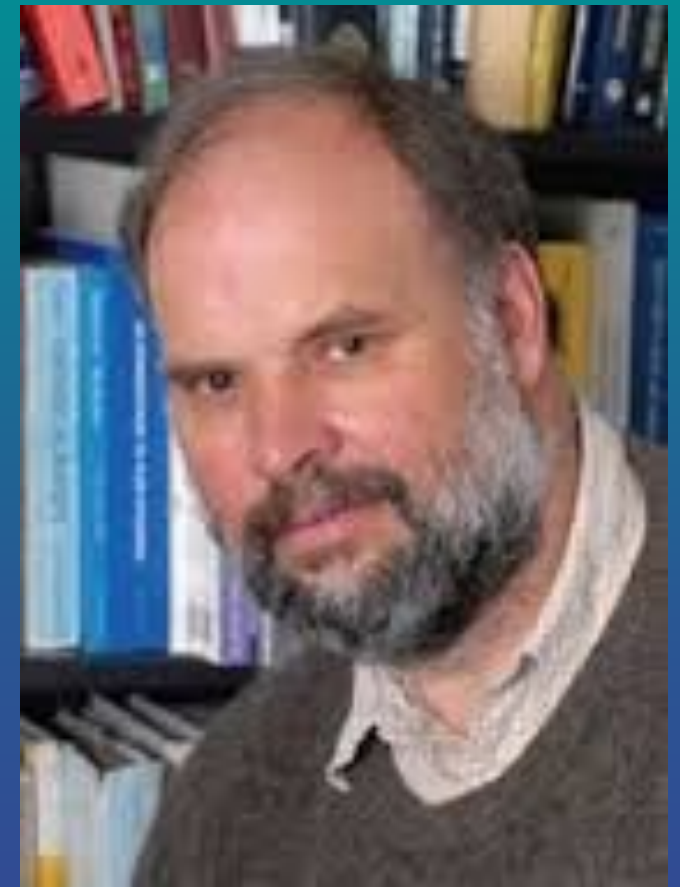
$K \neq K'$



$W(K) \neq W(K')$



Jones Polynomial



Kauffman Polynomial

Kauffman Bracket \sim Jones Polynomial



$= a$



$+ a^{-1}$





$= -a^2 - a^{-2}$


Kauffman Bracket \sim Jones Polynomial



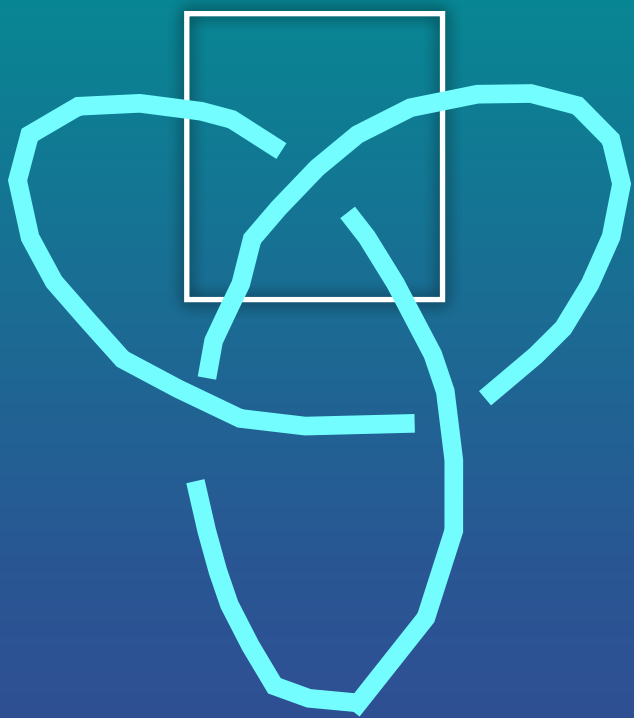
$= a$  $+ a^{-1}$ 

The diagram shows a positive crossing of two orange strands. To its right is an equals sign followed by the variable a . Further right are two parallel vertical orange strands. To their right is a plus sign followed by a^{-1} . Finally, to the right of a^{-1} are two parallel horizontal orange strands.



$= a^{-1}$  $+ a$ 

The diagram shows a negative crossing of two orange strands. To its right is an equals sign followed by a^{-1} . Further right are two parallel vertical orange strands. To their right is a plus sign followed by a . Finally, to the right of a are two parallel horizontal orange strands.



= a



+ a⁻¹



Opening each crossing



Two simpler diagrams

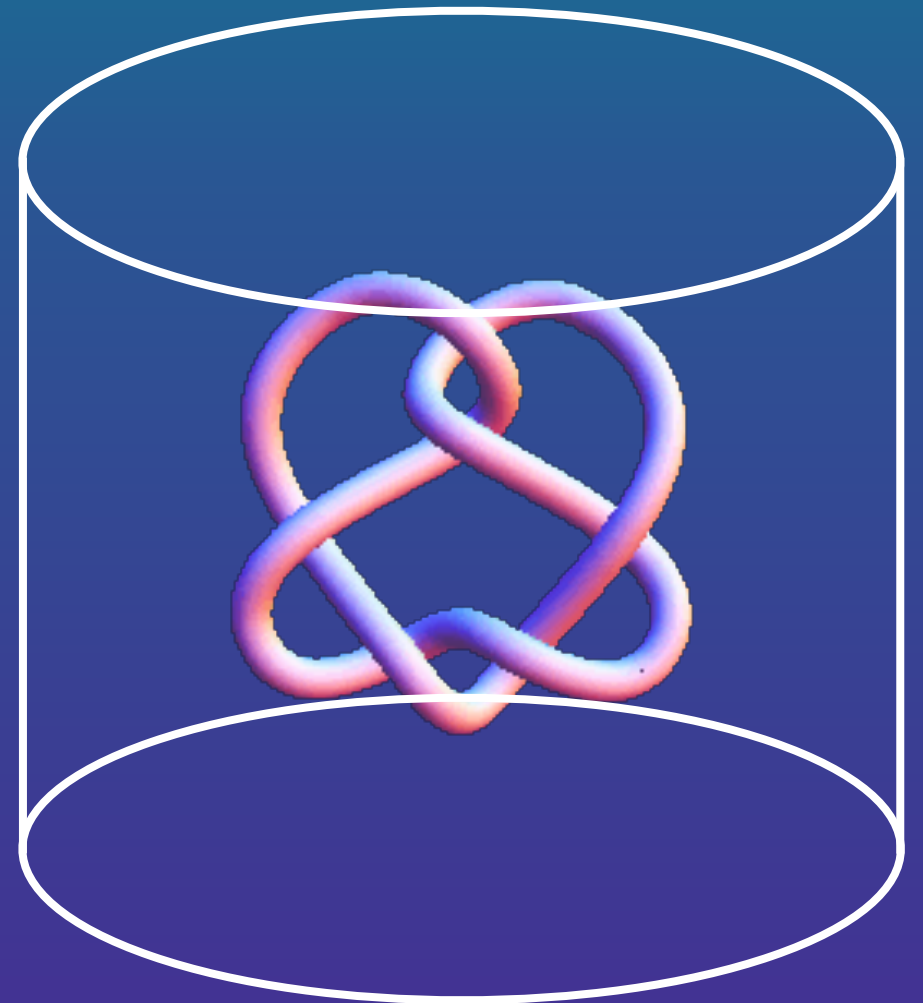
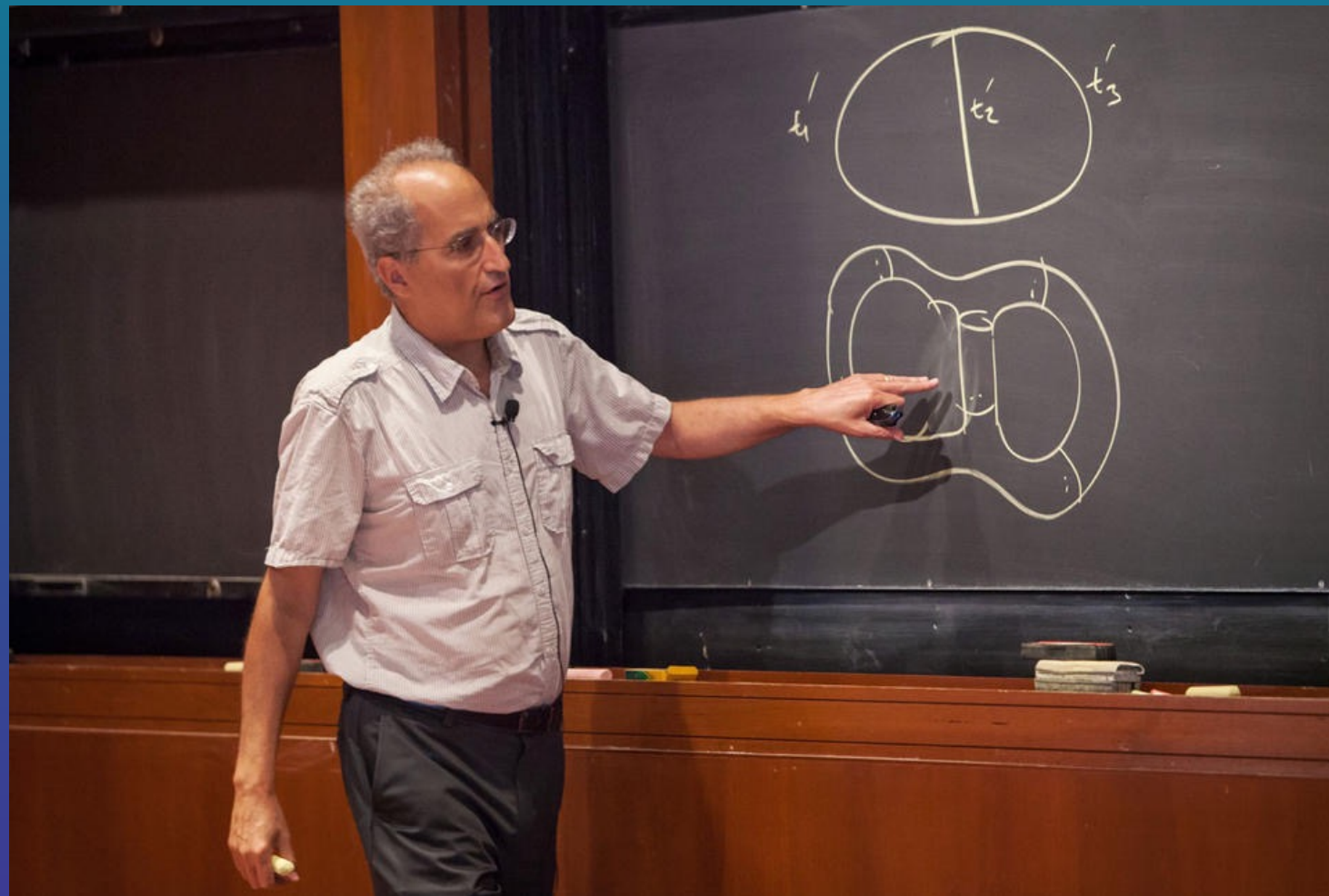
A computationally Hard Problem



2^N Crossings

Quantum Field Theory and the Jones Polynomial

$$Z = \int dA e^{iS}$$



$$\langle W_K \rangle = \int dA W_K[A] e^{iS}$$

$$S = \int d^3x \operatorname{Tr} \left[A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right]$$



?



Michael Freedman, Microsoft Station Q
University of California, Santa Barbara

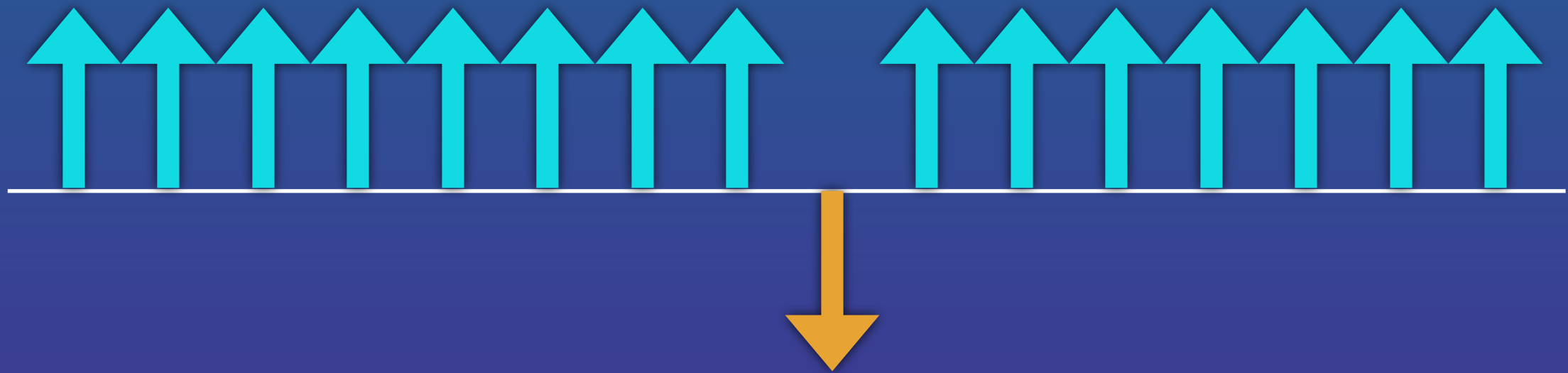
Alexi Kitaev,
Caltech

Topological Quantum Computation

Error Free Quantum Computation

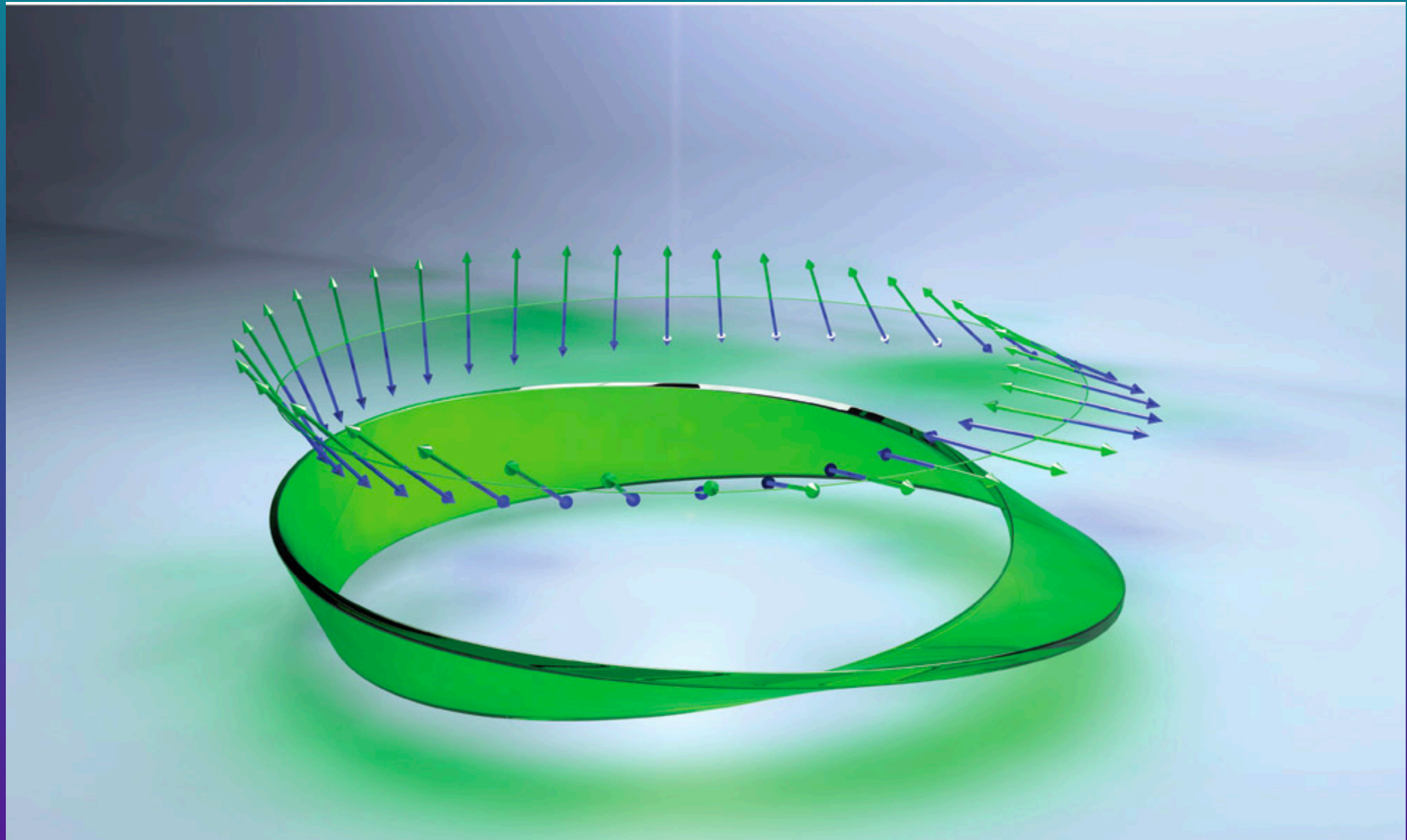
An excitation (a quasi particle) which is not topological

Spins in Ising Model



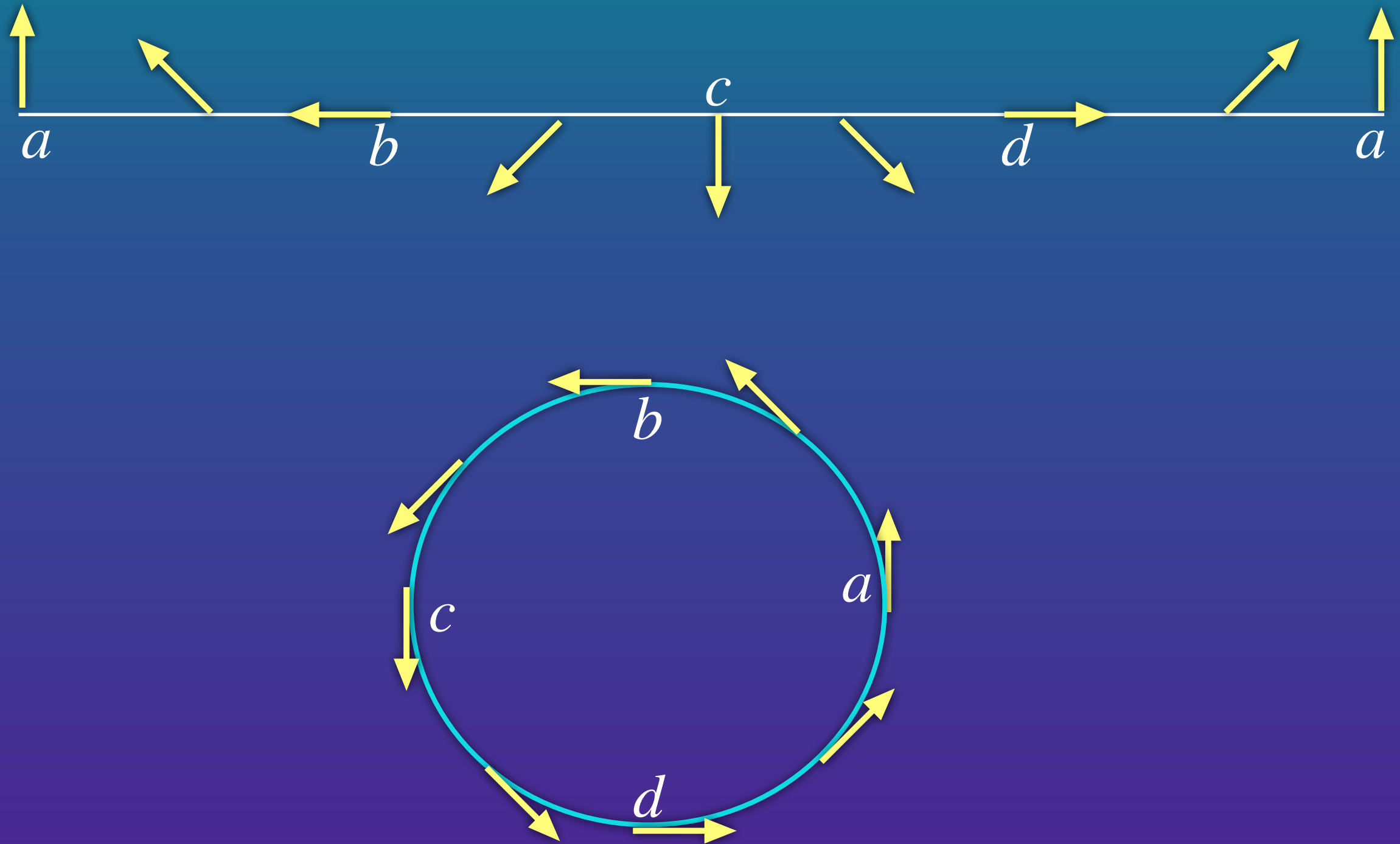
You flip the spin by a local operator and the excitation is removed!

An excitation (a quasi particle) which is topological

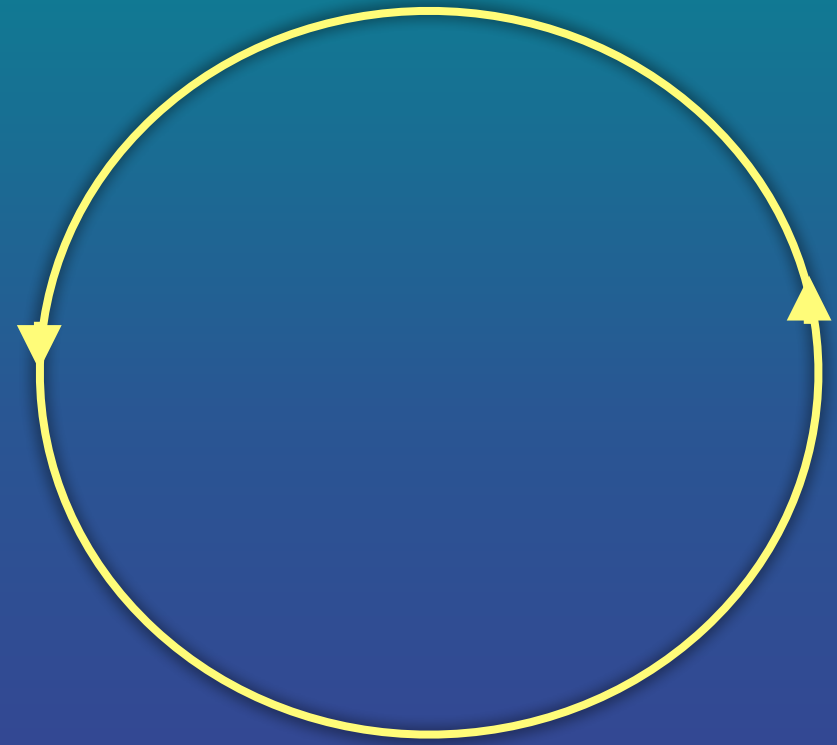
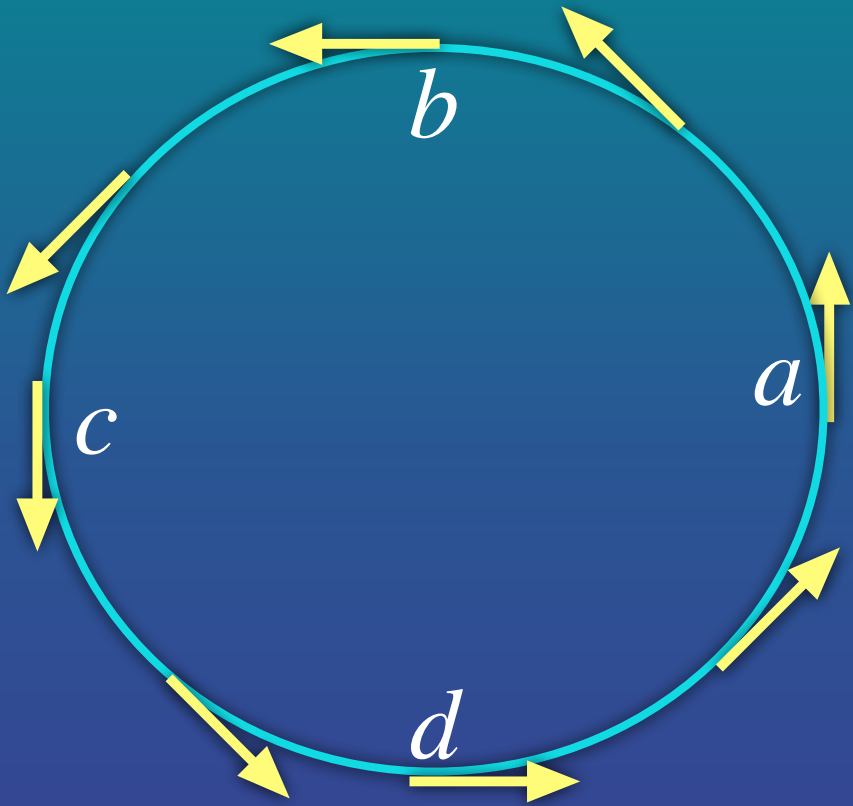


You cannot remove the excitation by local operations. flip the spin by a local operator and the excitation is removed!

An excitation (a quasi-particle) which is topological.



Topology protects the excitation.

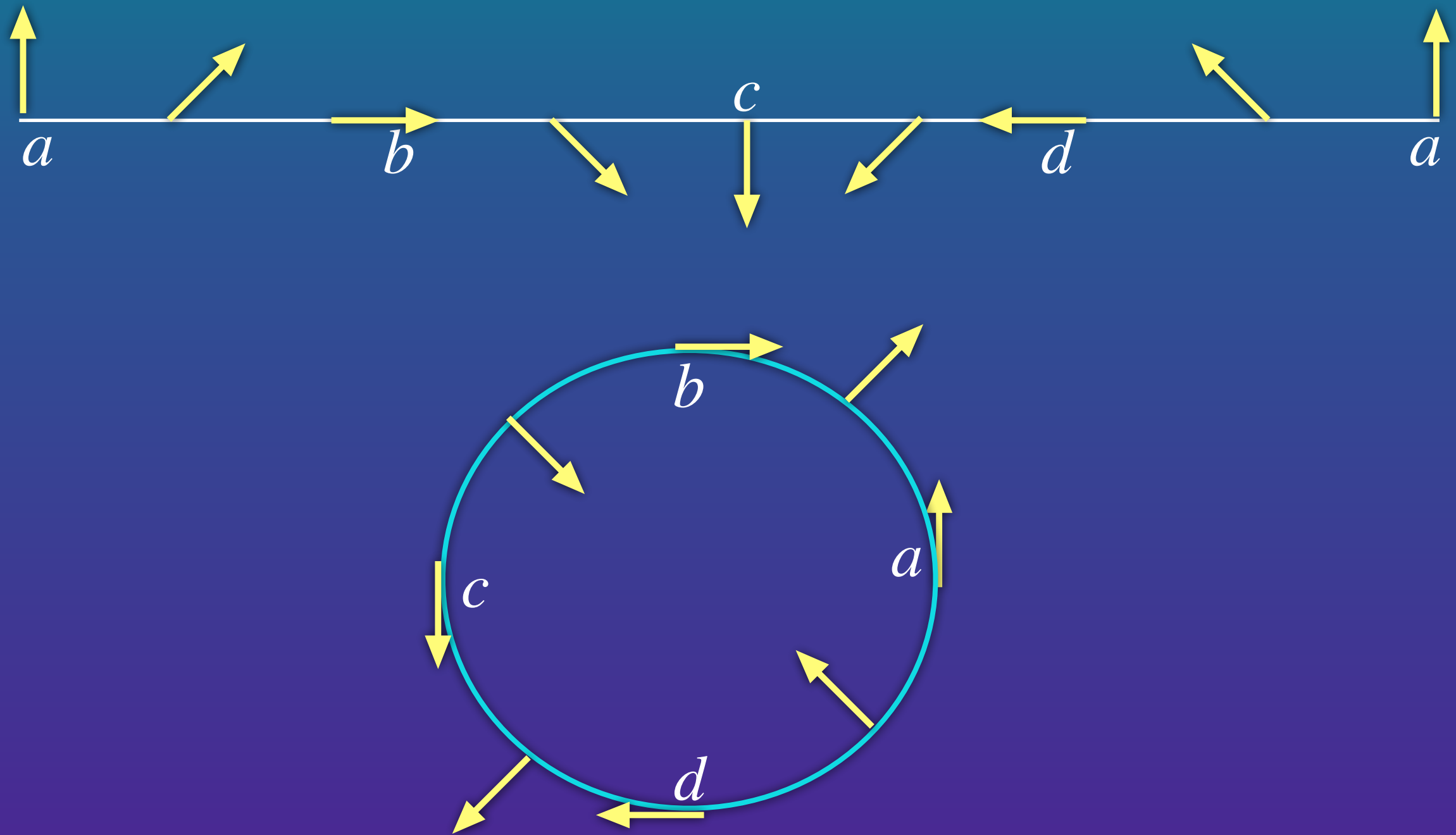


$$\phi : S^1 \longrightarrow S^1$$

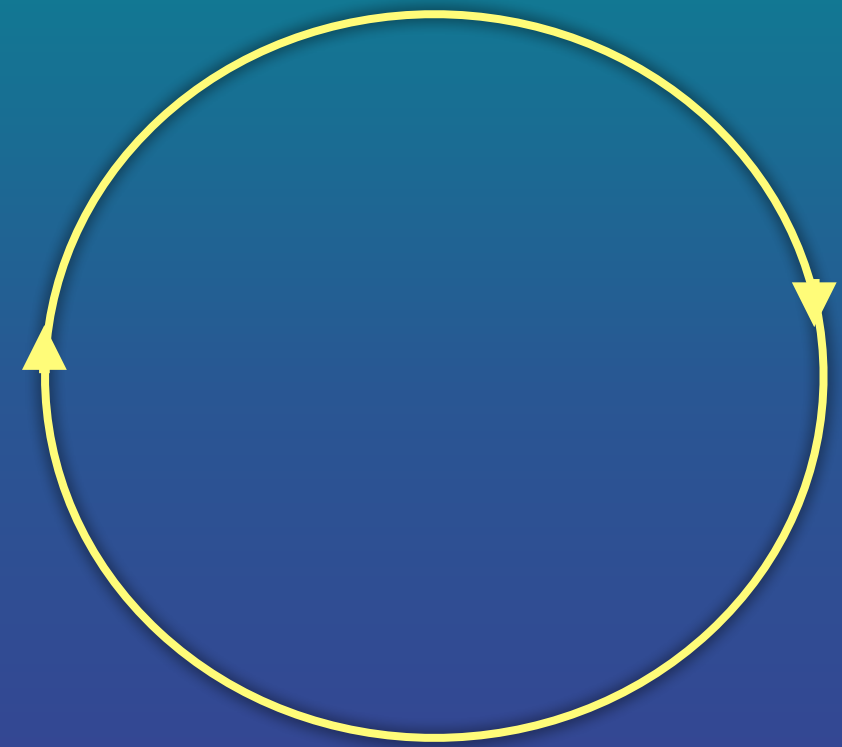
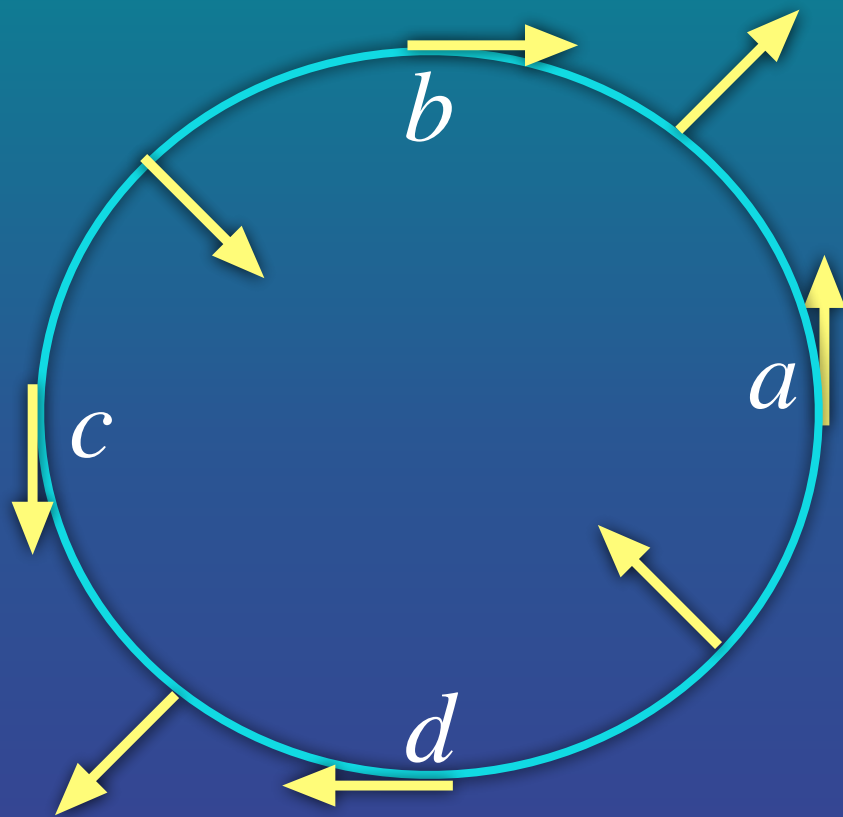
$$\phi : \text{Space} \longrightarrow \text{Spin}$$

Winding number = $q = 1$

An excitation (a quasi-particle) which is topological.



Topology protects the excitation.

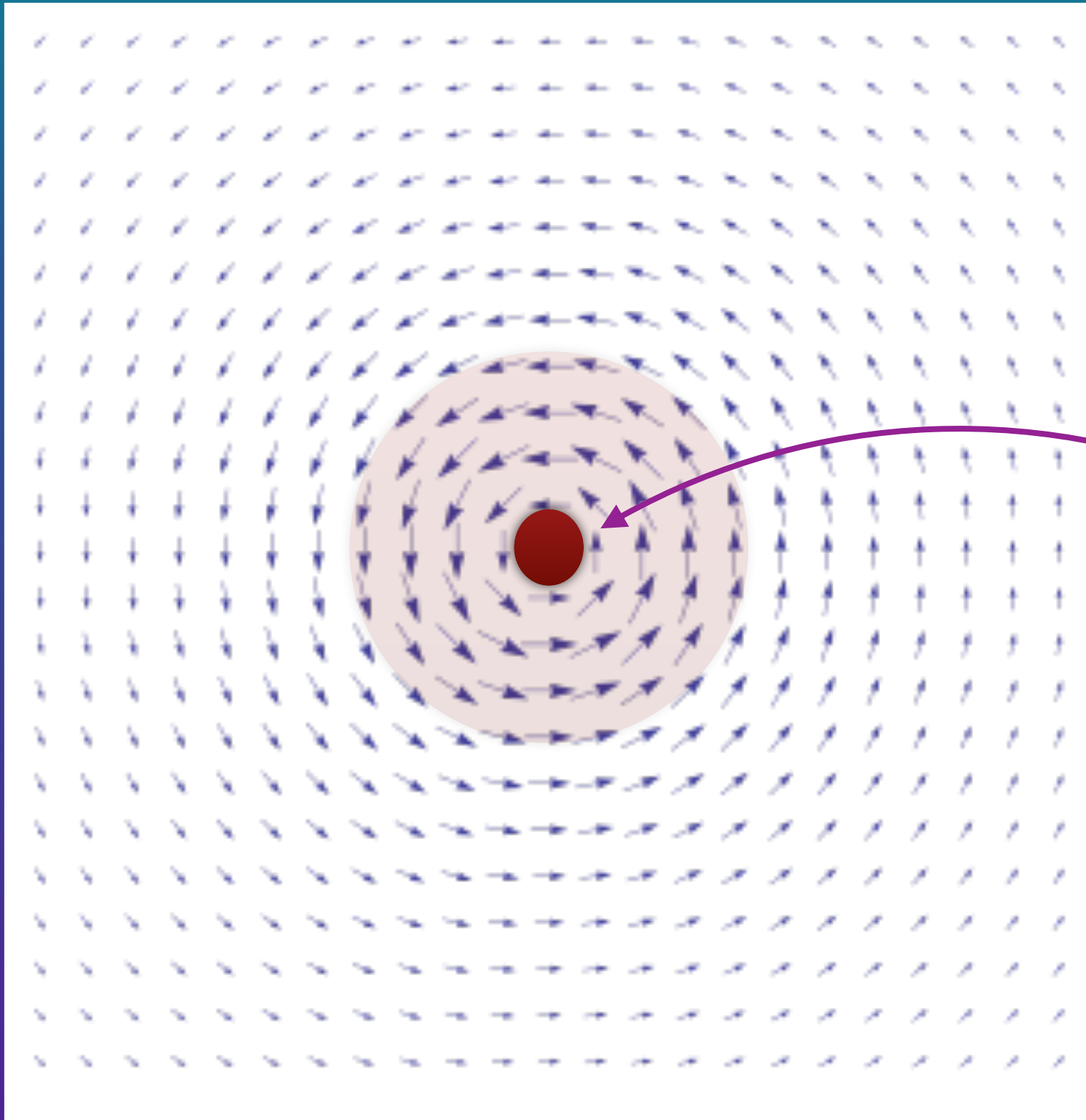


$$\phi : S^1 \longrightarrow S^1$$

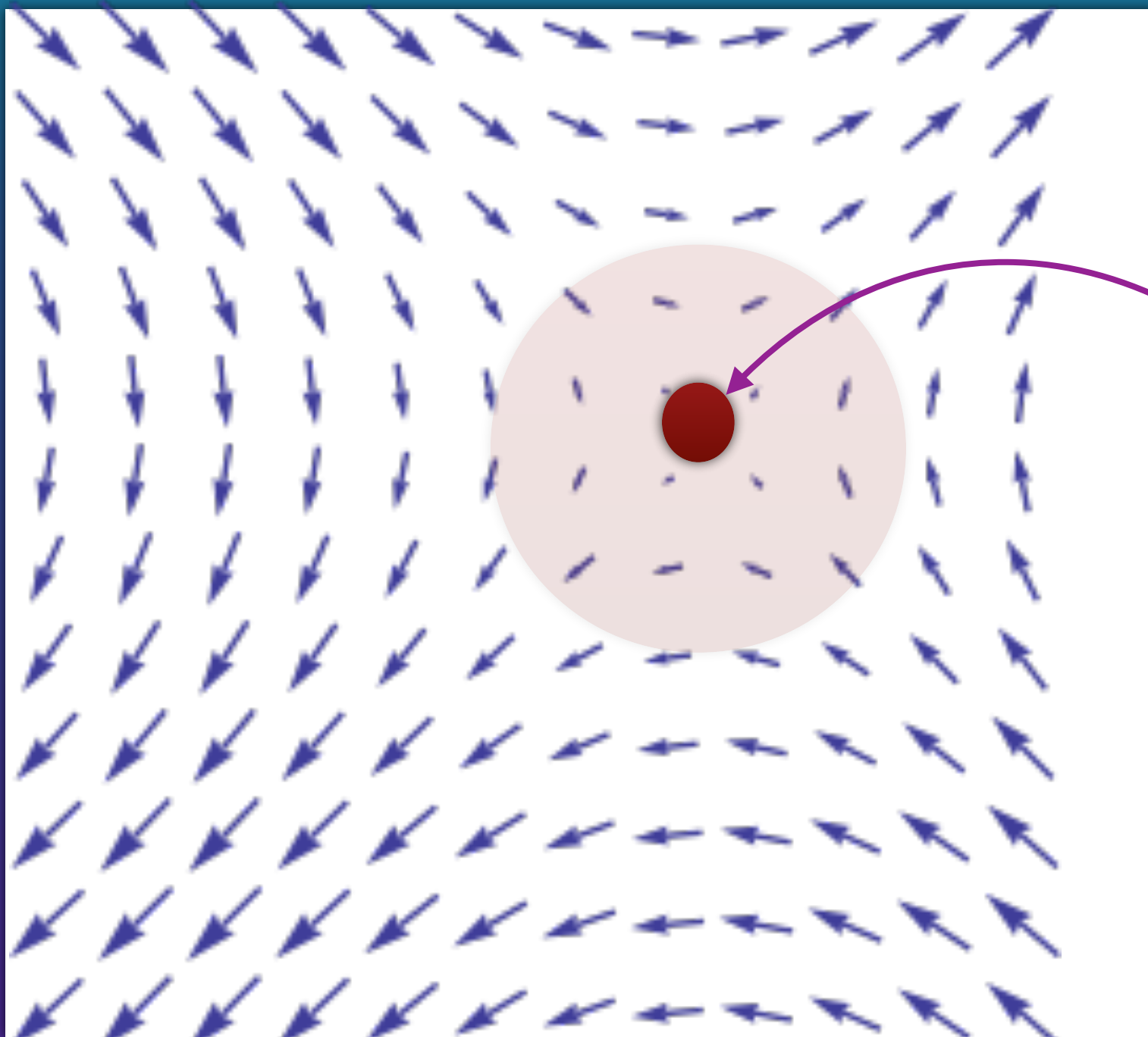
$$\phi : \text{Space} \longrightarrow \text{Spin}$$

Winding number = $q = -1$

Quasi particles have charge!



Anyon with charge $q=1$

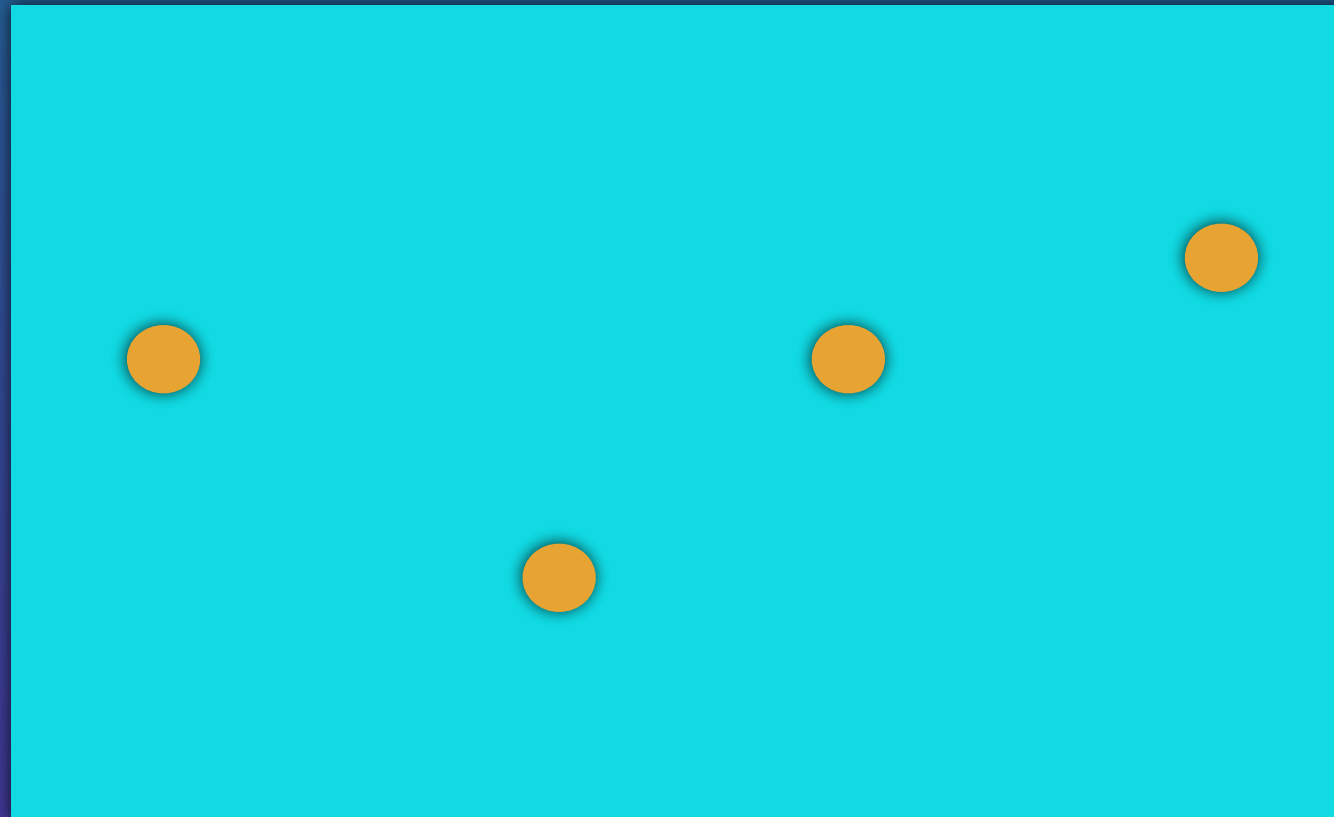


Anyon with charge $q=1$

The basic principles

- 1- There are many body systems whose ground states have topological charge

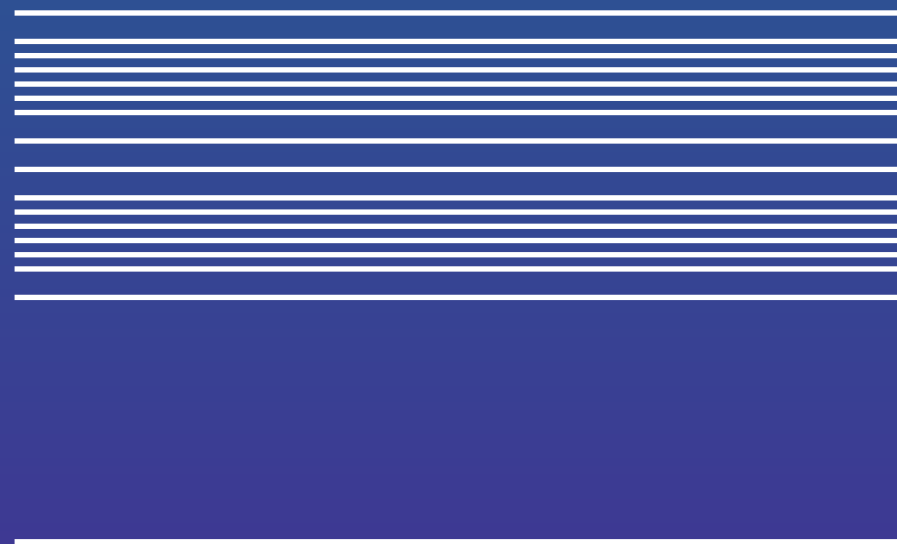
2-The degeneracy of the ground state depends on the number of these particles.



The higher the number of charges \longrightarrow The higher the degeneracy

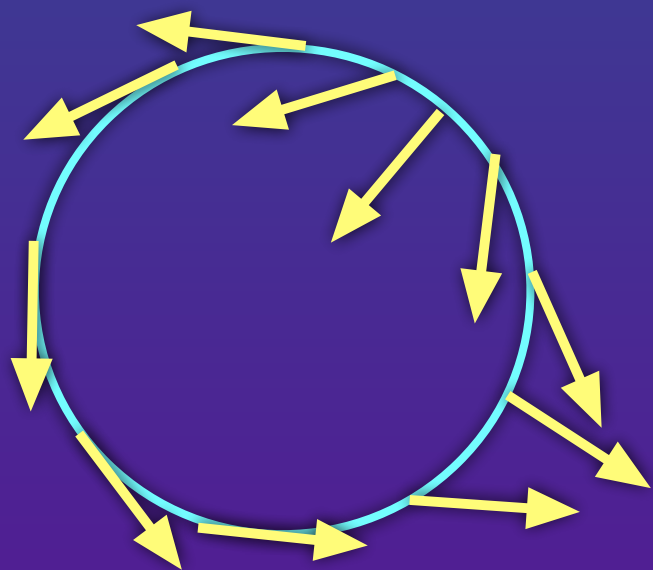
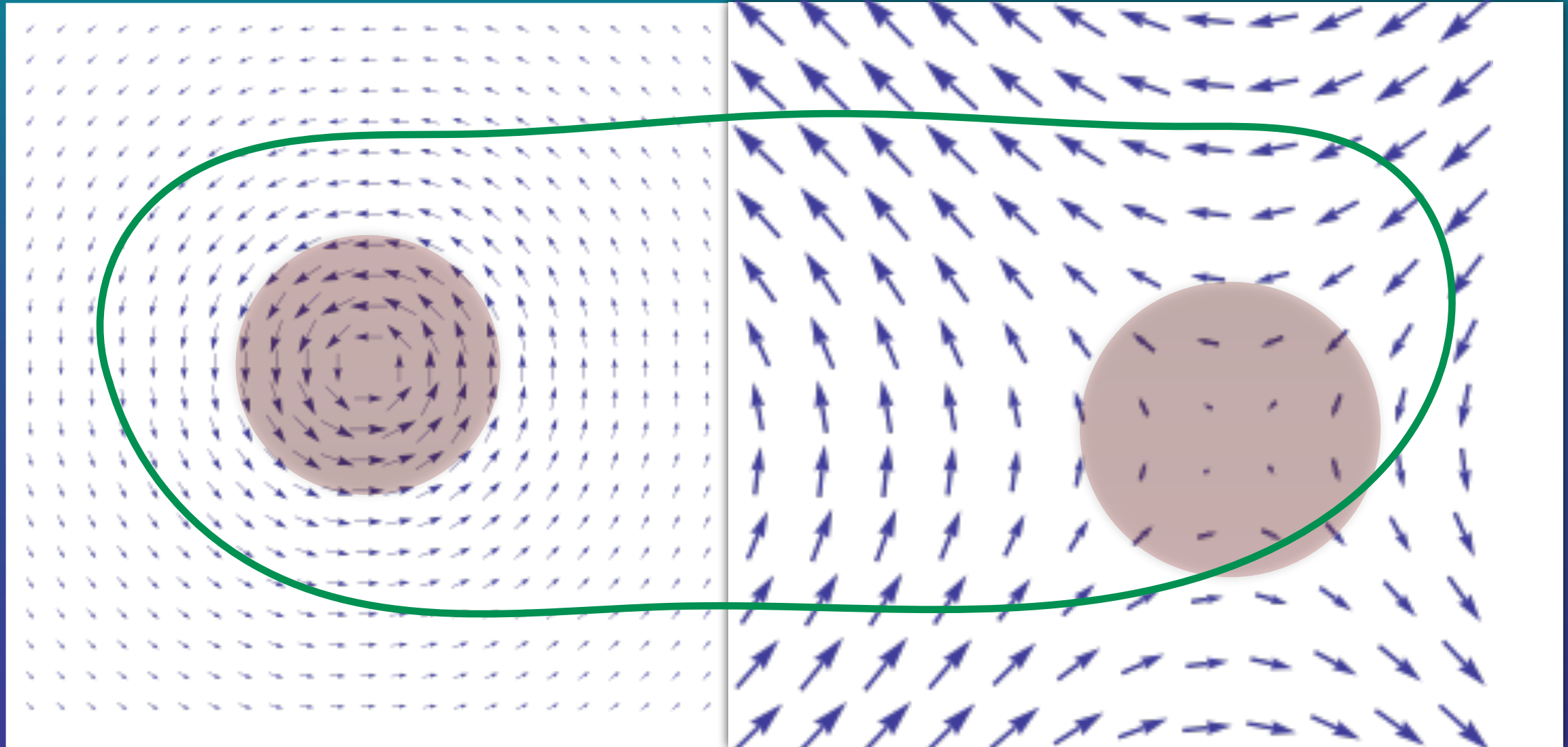
$$|g_1\rangle, |g_2\rangle, |g_3\rangle, \dots, |g_N\rangle$$

3-The system has a gap.



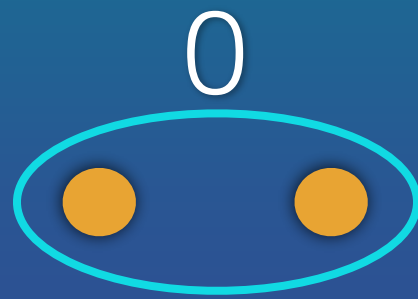
We always stay in the ground space.

Combination of topological charges

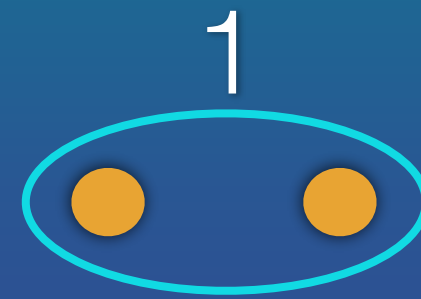


$$q = 1 + (-1) = 0$$

Example: **Fibonacci Anyons** 0 and 1



$$1 \times 1 = 0$$

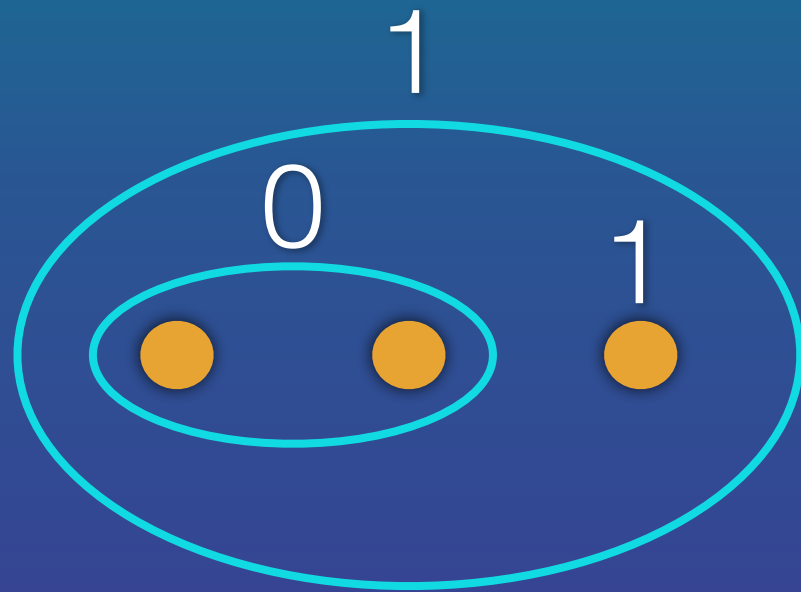


$$1 \times 1 = 1$$

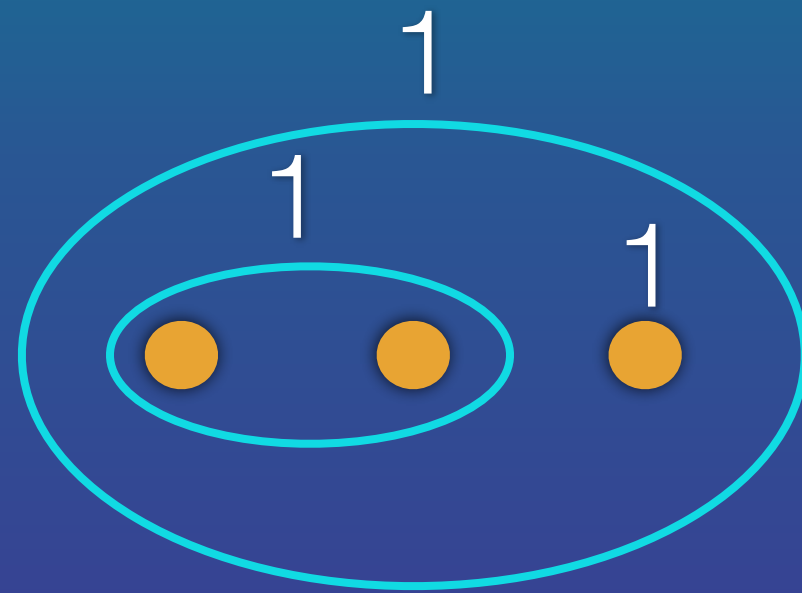
$$1 \times 1 = 0 + 1$$

By two Fibonacci Anyons, we cannot make a qubit since their total charge will be different and so we cannot make linear superposition of them

By three Fibonacci Anyons, we can make a qubit

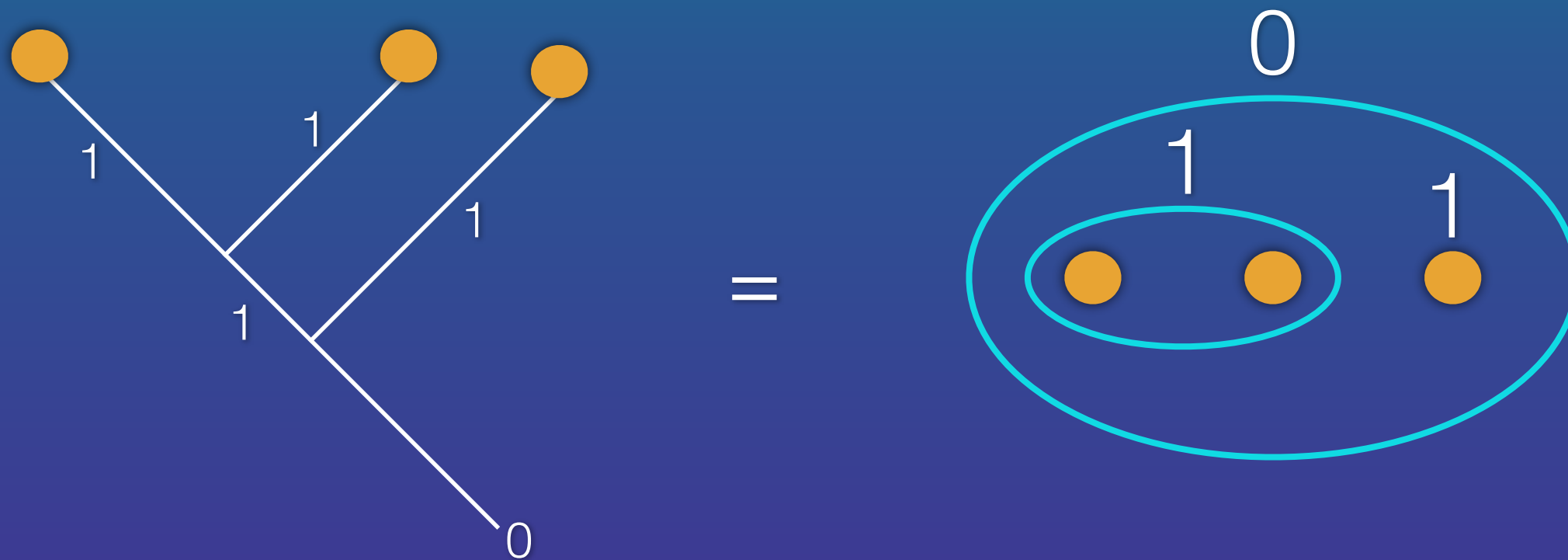


$|0\rangle$



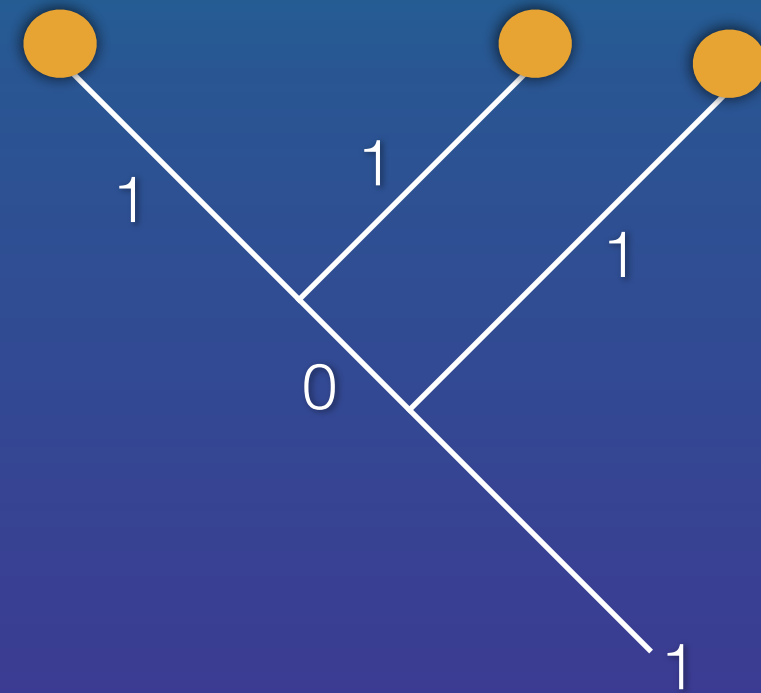
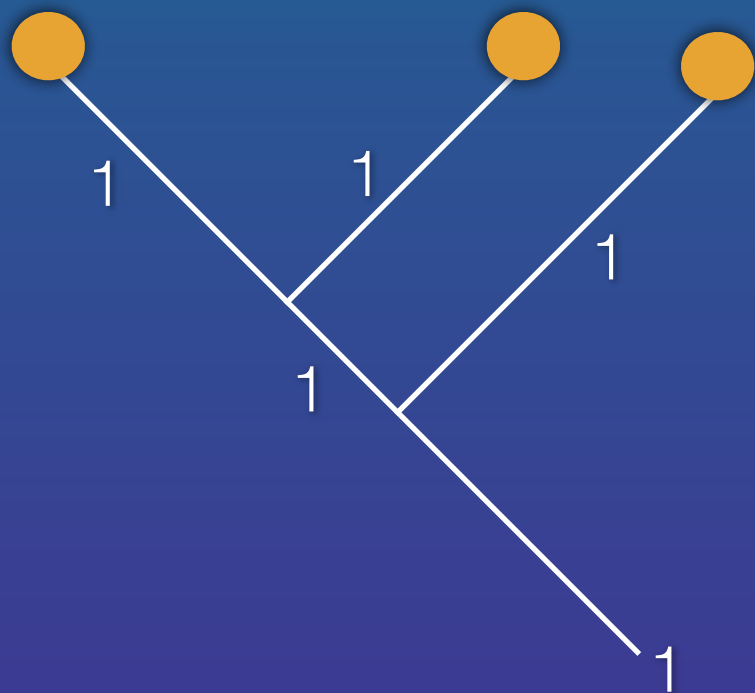
$|1\rangle$

With three Fibonacci anyons with total charge zero, we cannot make a qubit.



Here is there is no degeneracy.

With three Fibonacci anyons with total charge one, we can make a qubit.



Here is there is degeneracy.

The ground state with n Fibonacci anyons with total charge 1, has a degeneracy given by

$$\phi^n$$

$$\phi = \frac{\sqrt{5} + 1}{2}$$

ϕ = quantum dimension

$$1^n = a_n 0 + b_n 1$$

$$1^{n+1} = a_n 0 \times 1 + b_n 1 \times 1$$

$$= a_n 1 + b_n (0 + 1)$$

$$= b_n 0 + (a_n + b_n) 1$$

$$a_{n+1} = b_n$$

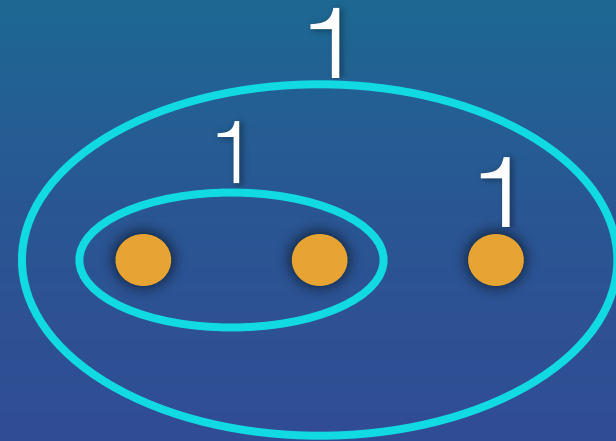
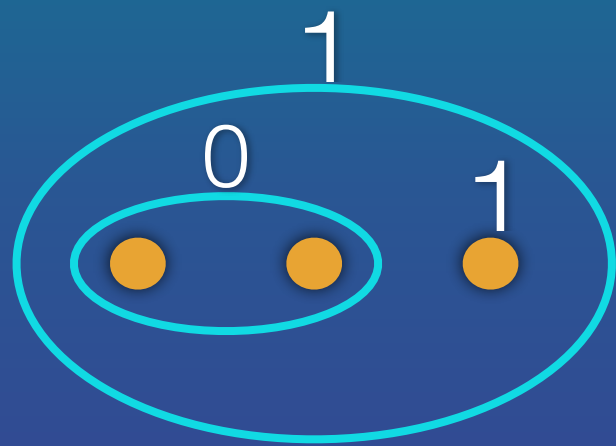
$$b_{n+1} = a_n + b_n$$

$$\begin{pmatrix} a_{n+1} \\ b_{n+1} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} a_n \\ b_n \end{pmatrix}$$

$$\phi_{\pm} = \frac{1 \pm \sqrt{5}}{2}$$

$$b_n \sim \phi_+^n$$

Can we really think of these



as qubits

$|0\rangle$

and

$|1\rangle$

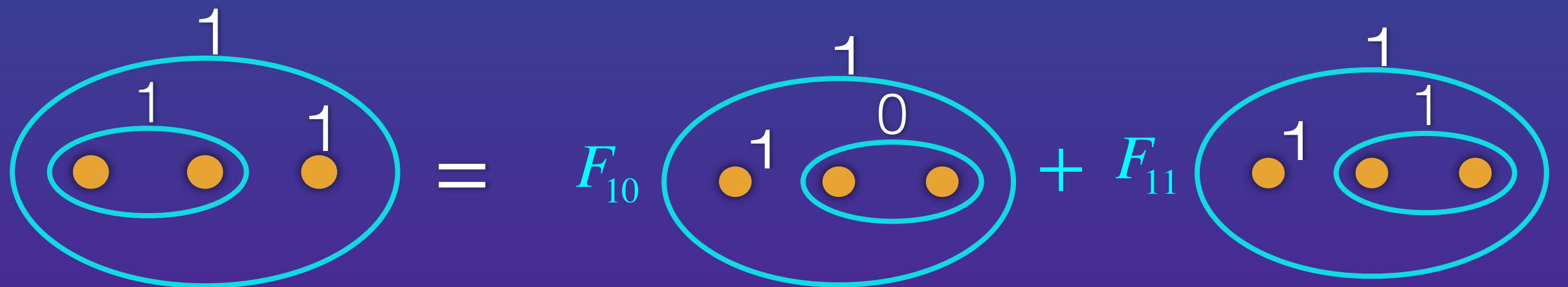
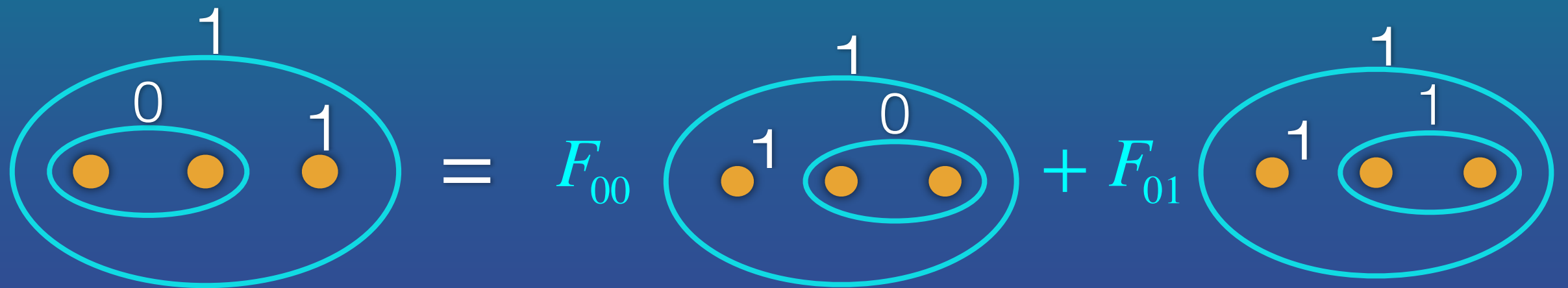
?

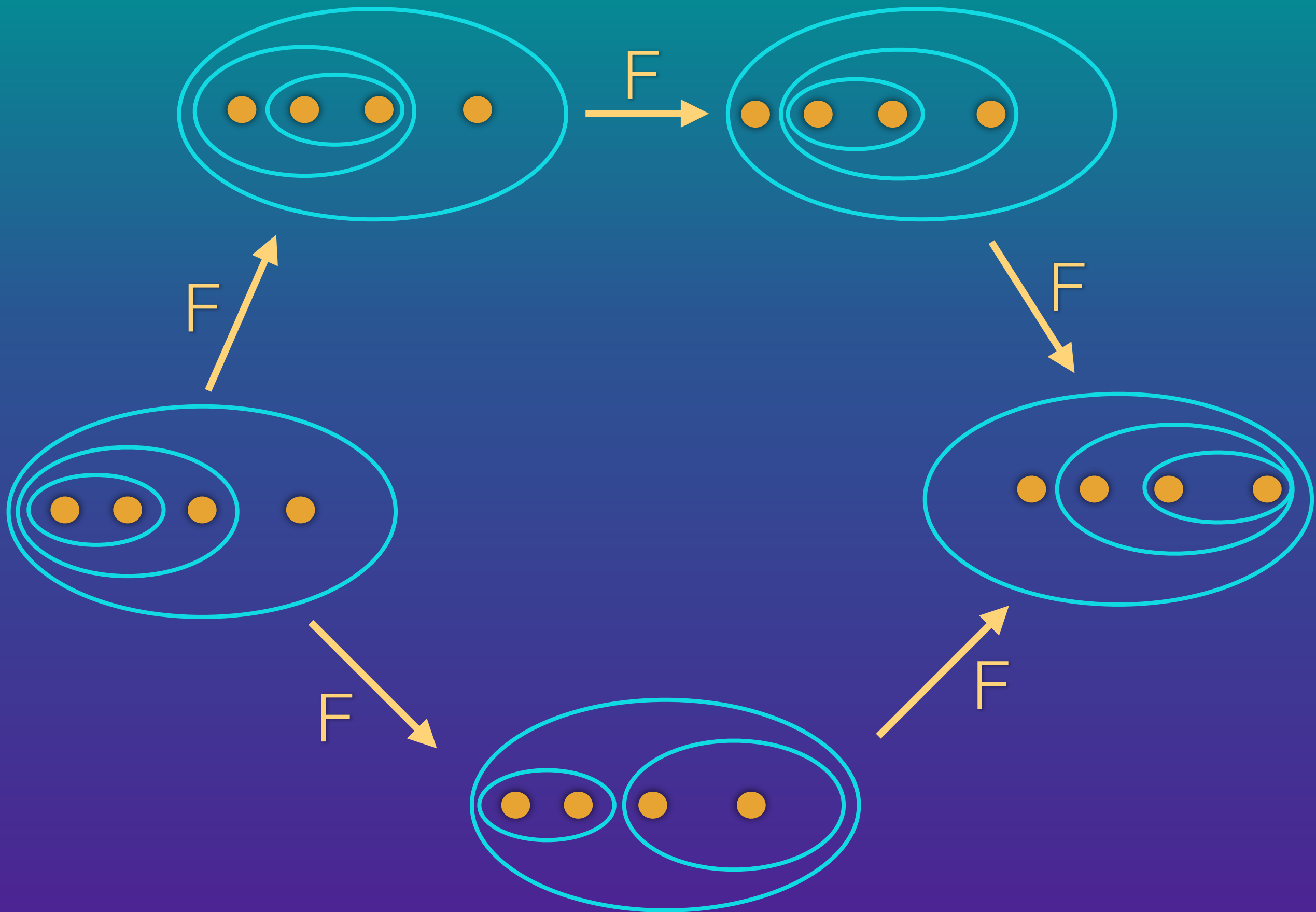
For example, what are the following?

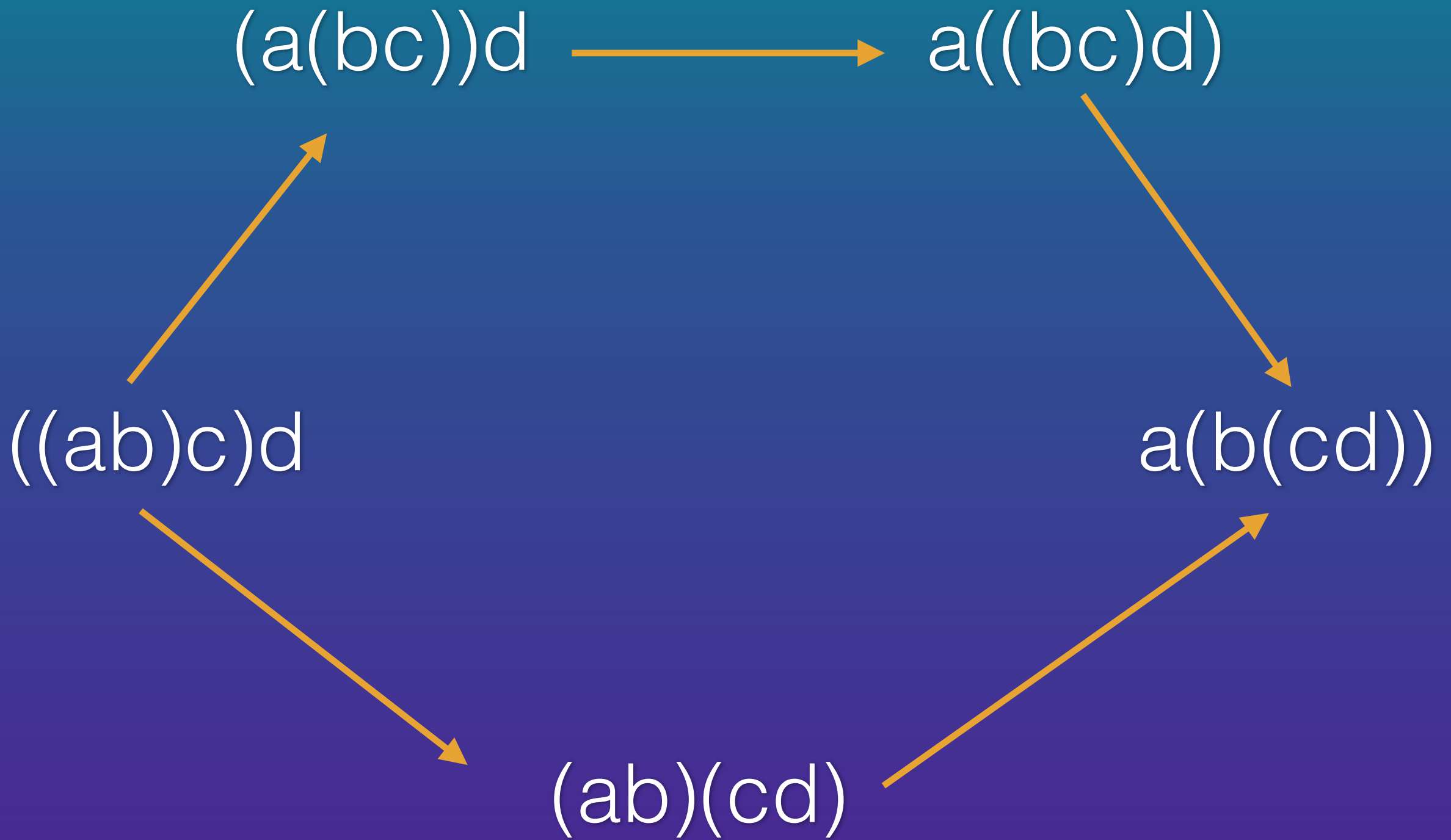


Answer: They should be linear combinations of the previous ones.

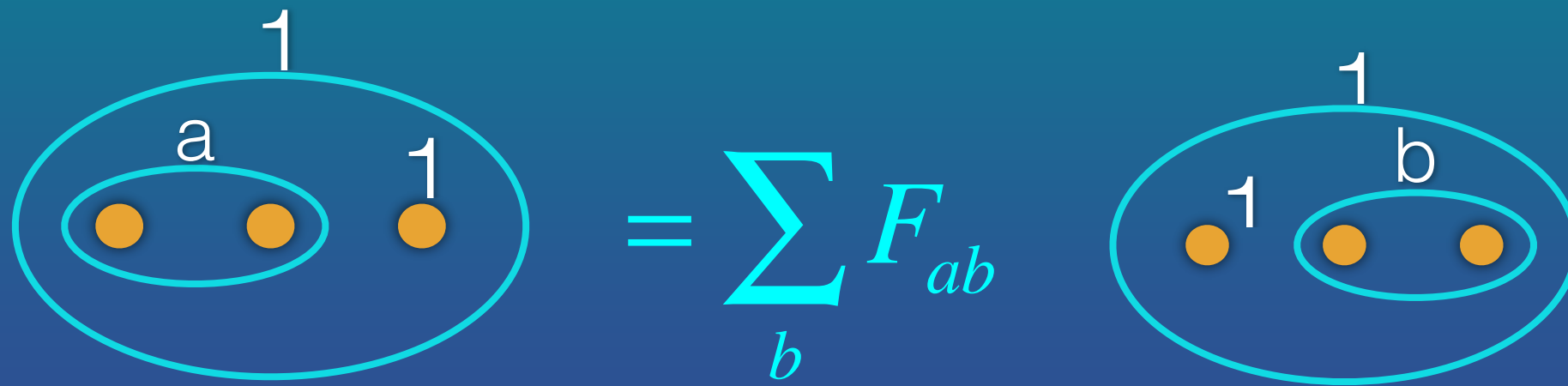
Consistency Conditions







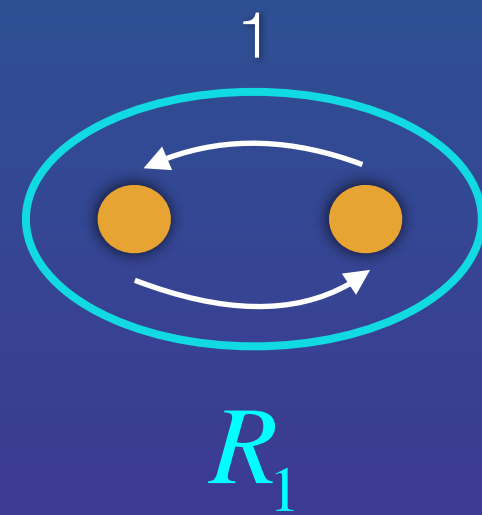
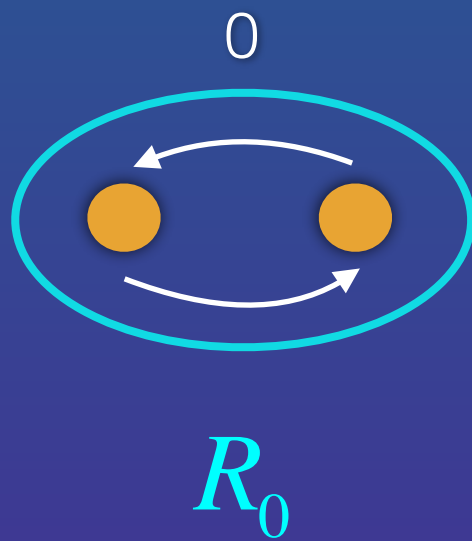
5-Consistency Relations?



$$F = \begin{pmatrix} \phi & \sqrt{\phi} \\ \sqrt{\phi} & -\phi \end{pmatrix}$$

$$\phi = \frac{\sqrt{5} - 1}{2}$$

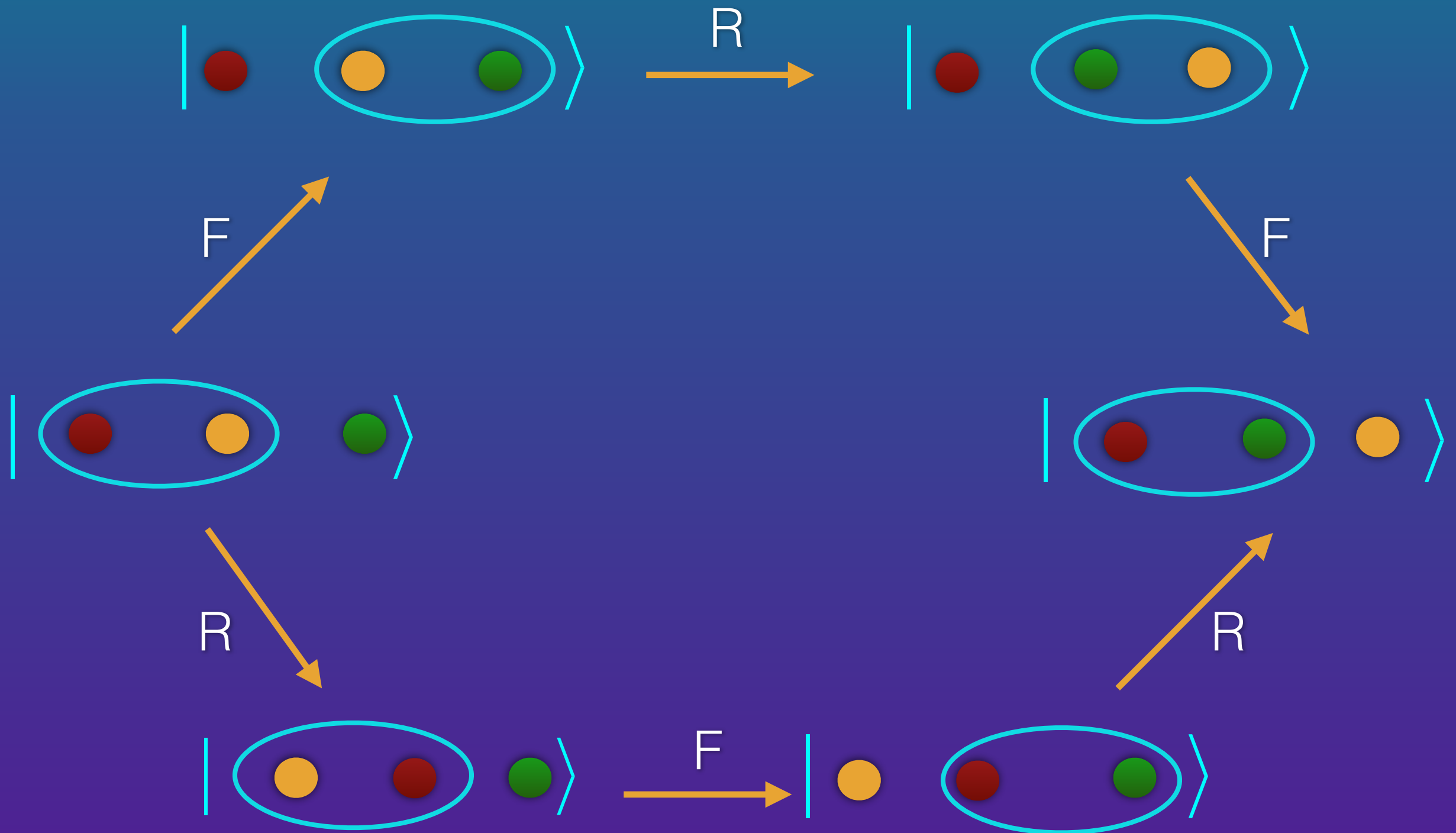
Braiding





We can do this in two different ways.

The charges obey certain fusion and braiding rules.

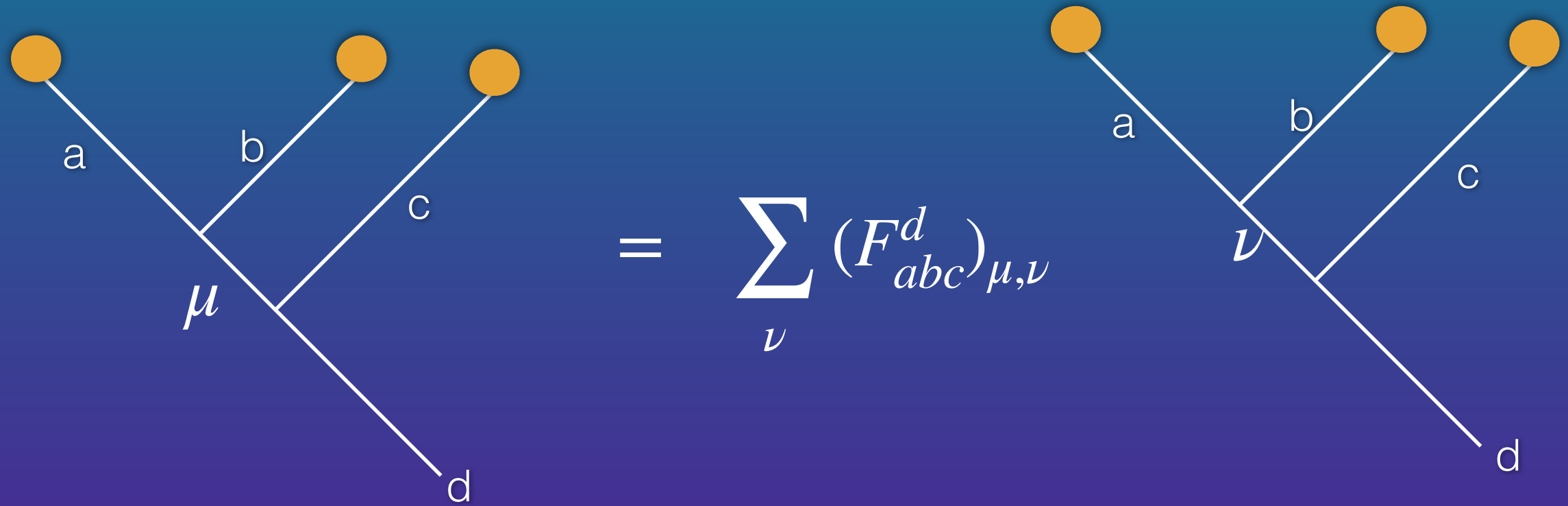


The braiding matrix

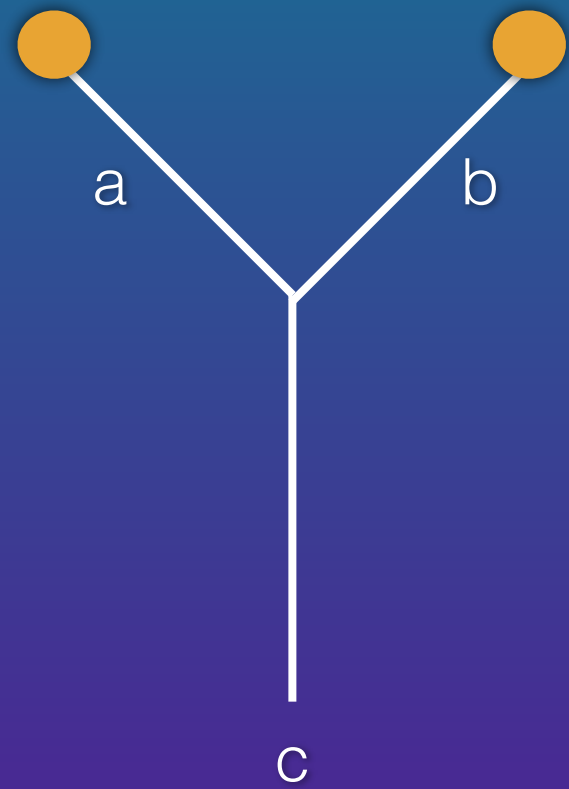


$$R = \begin{pmatrix} R_0 & \\ & R_1 \end{pmatrix} = \begin{pmatrix} e^{\frac{4\pi i}{5}} & \\ & -e^{\frac{2\pi i}{5}} \end{pmatrix}$$

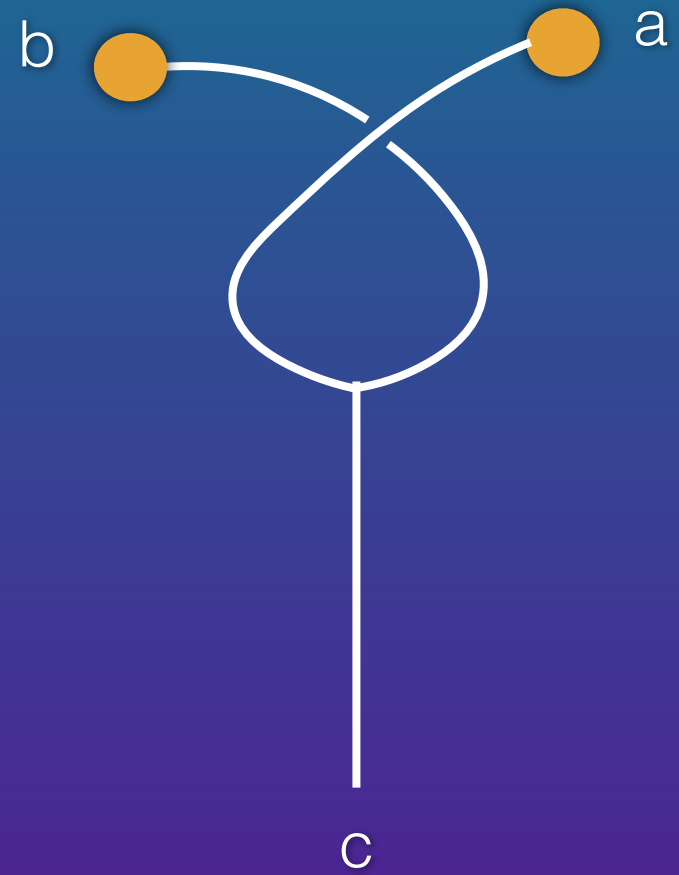
It is not so simple!



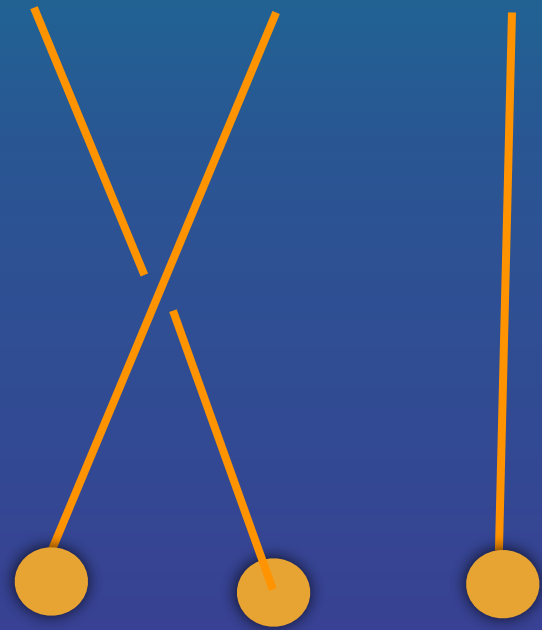
It is not so simple!



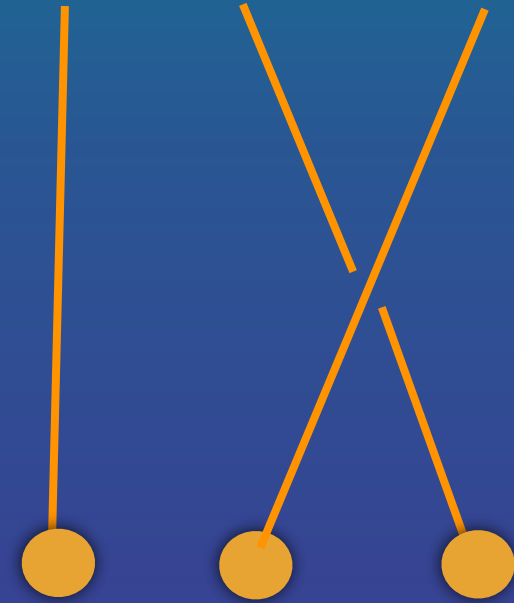
$$= R_{ab}^c$$



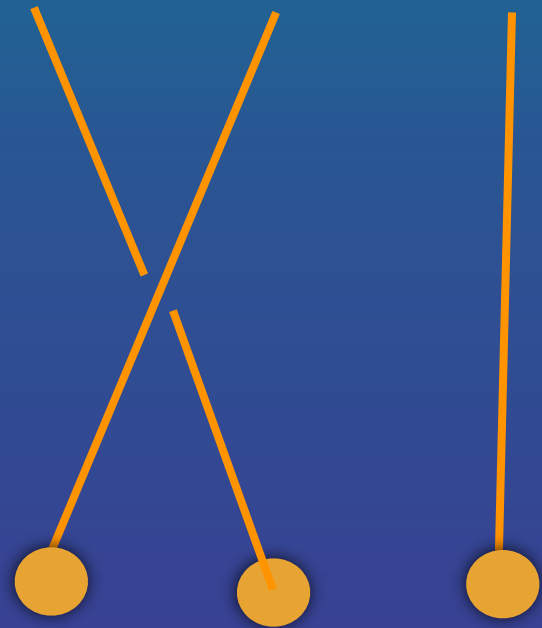
Quantum Gates



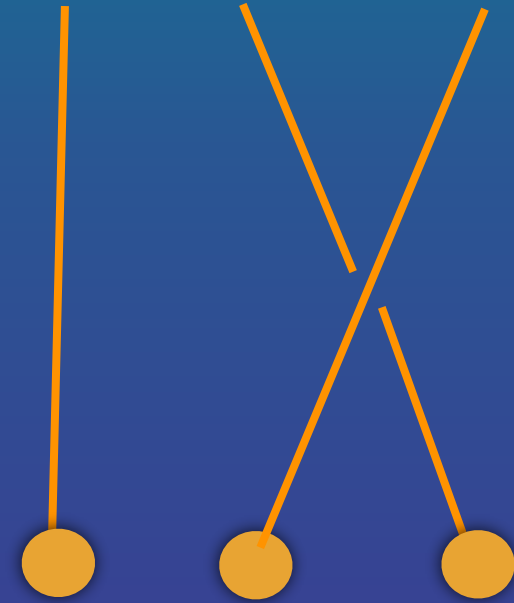
$$\sigma_1 = R$$



$$\sigma_2 = FRF$$



$$\sigma_1 = R = R_z\left(\frac{3\pi}{10}\right)$$



$$\sigma_2 = FRF = R_n(\theta)$$



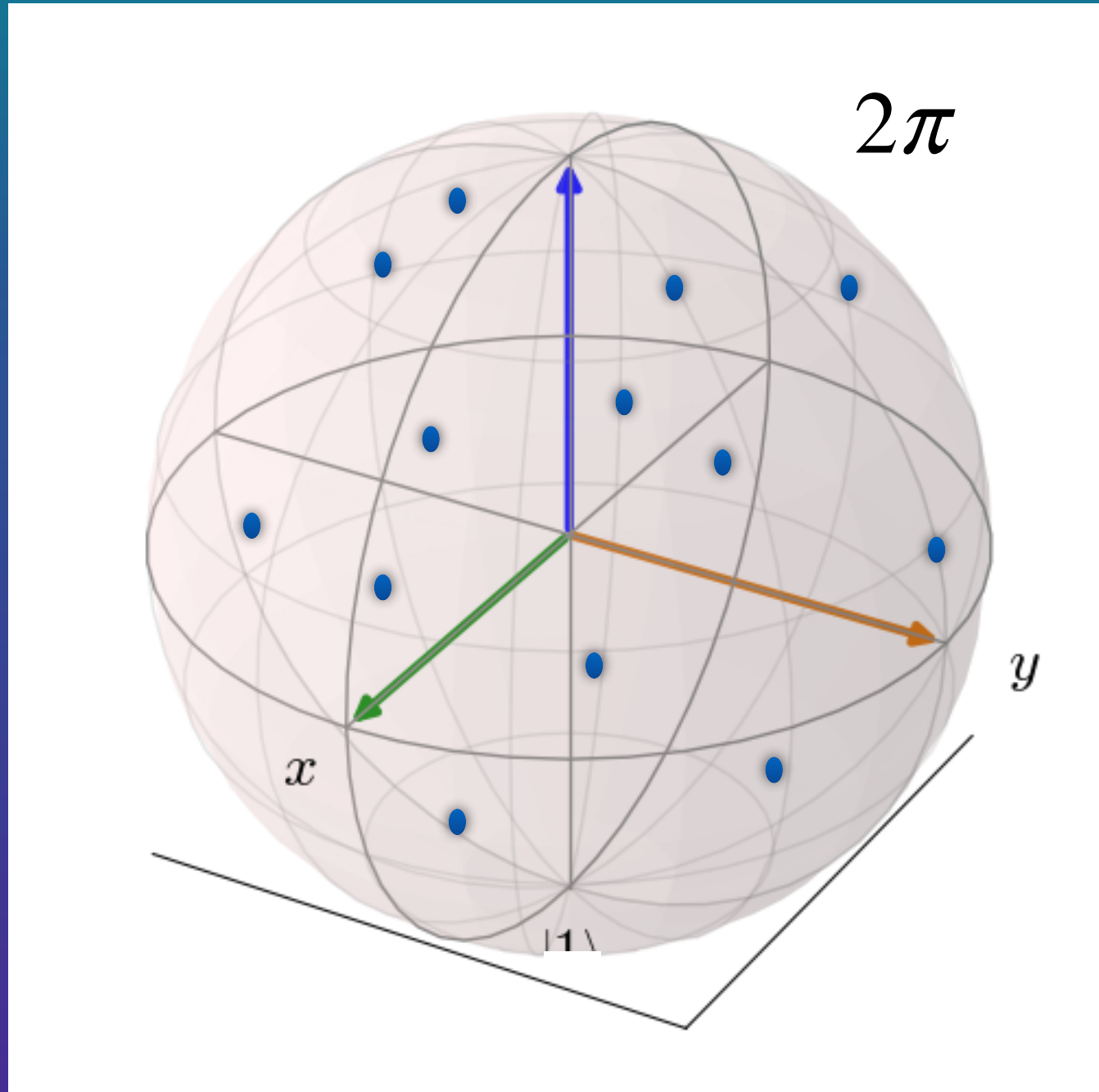
$$R_z\left(\frac{6\pi}{10}\right)$$

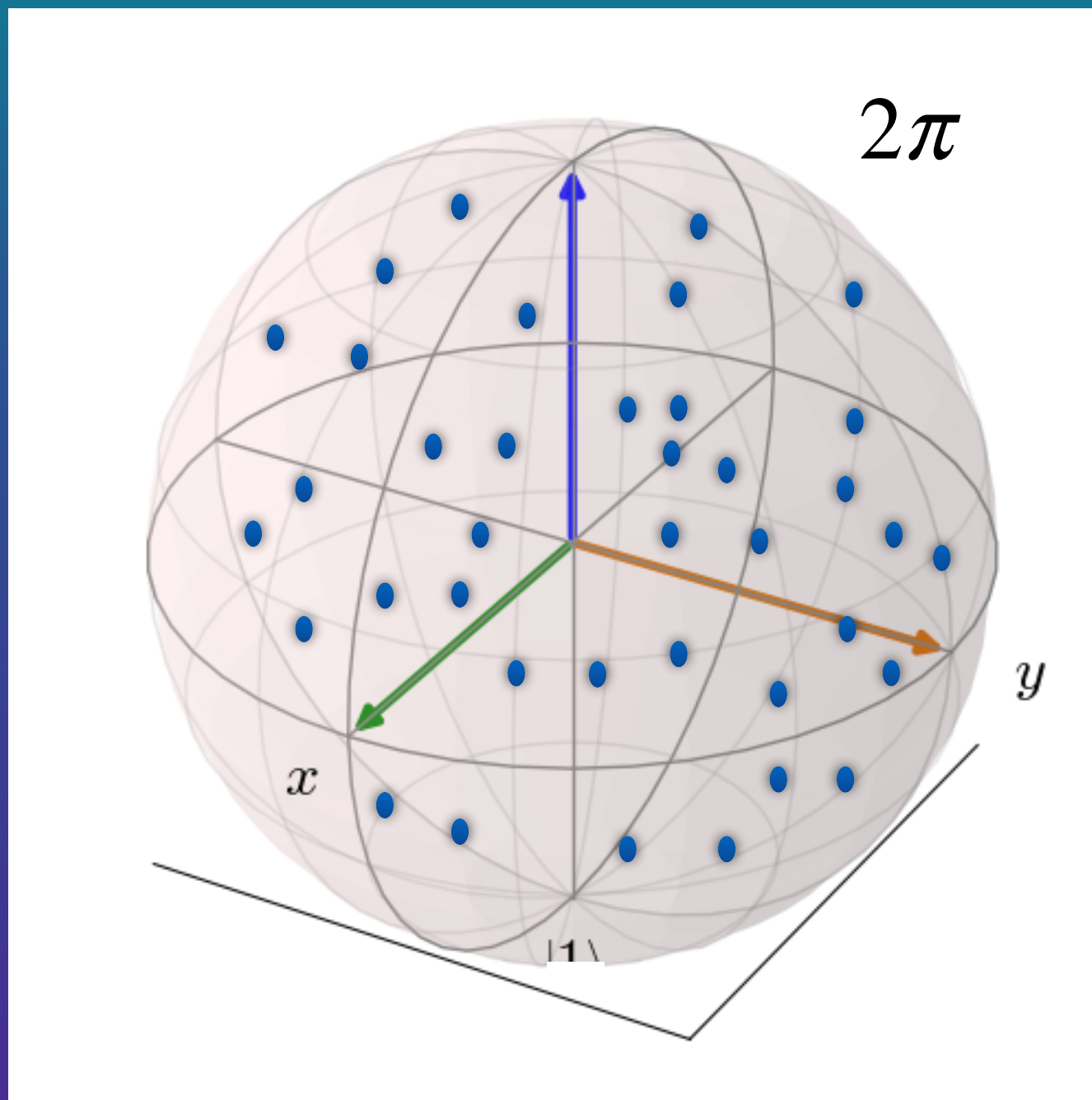


$$R_n(2\theta)$$



$$\sigma_1 \sigma_2 \sigma_2 \sigma_1 \sigma_1$$

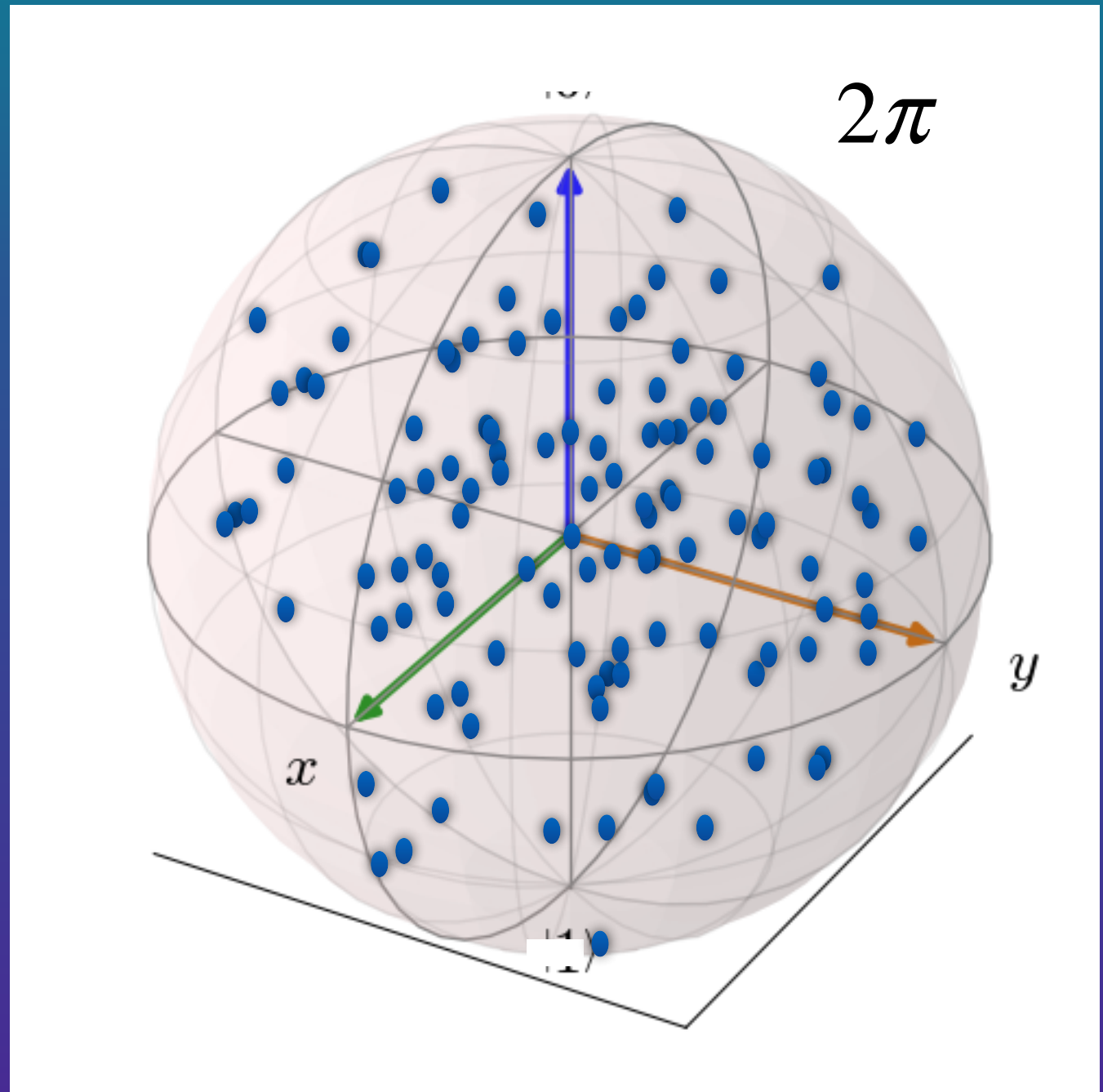




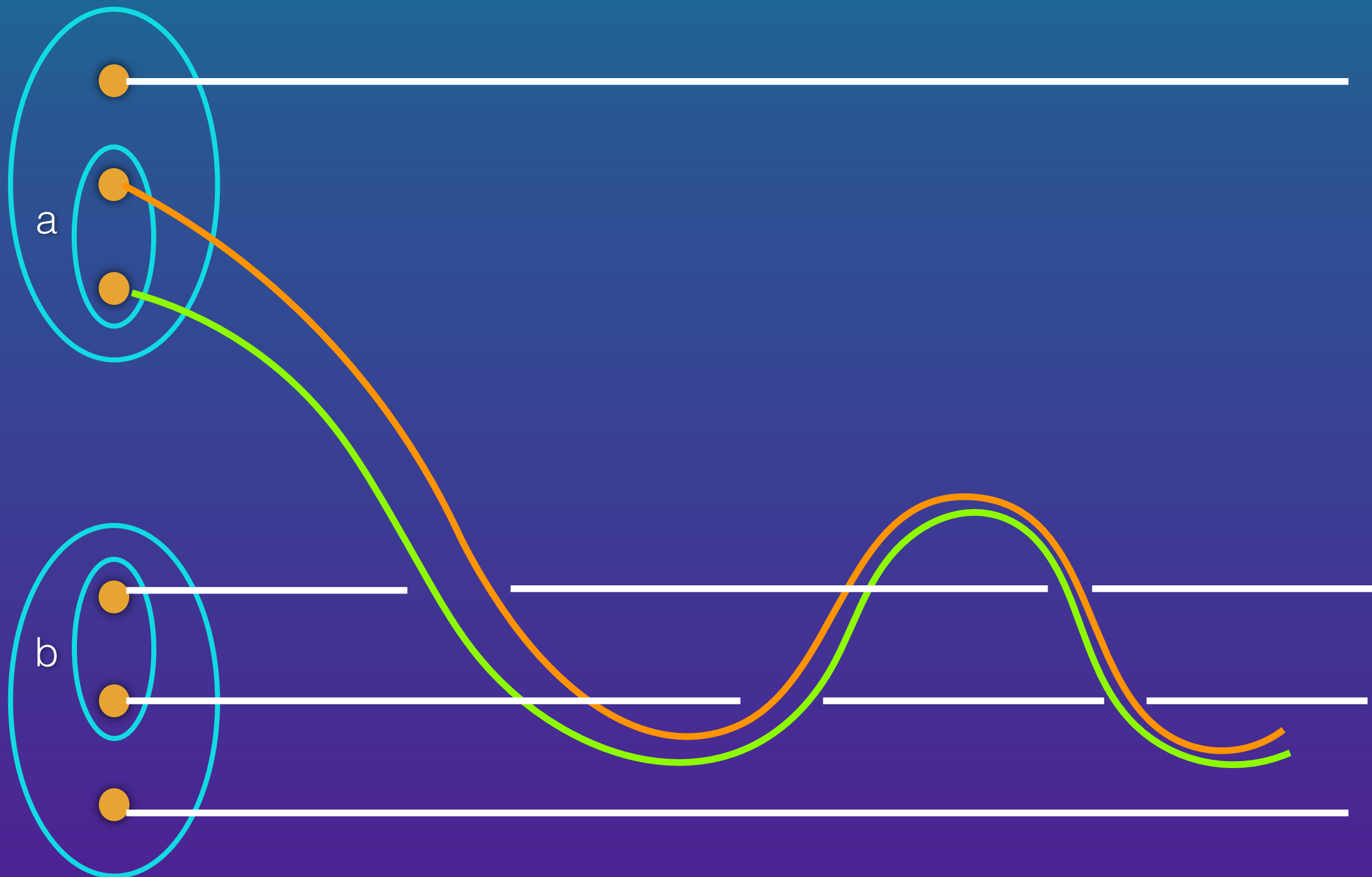
$\sigma_1 \sigma_2 \sigma_2 \sigma_1 \sigma_1 \sigma_2 \sigma_2$



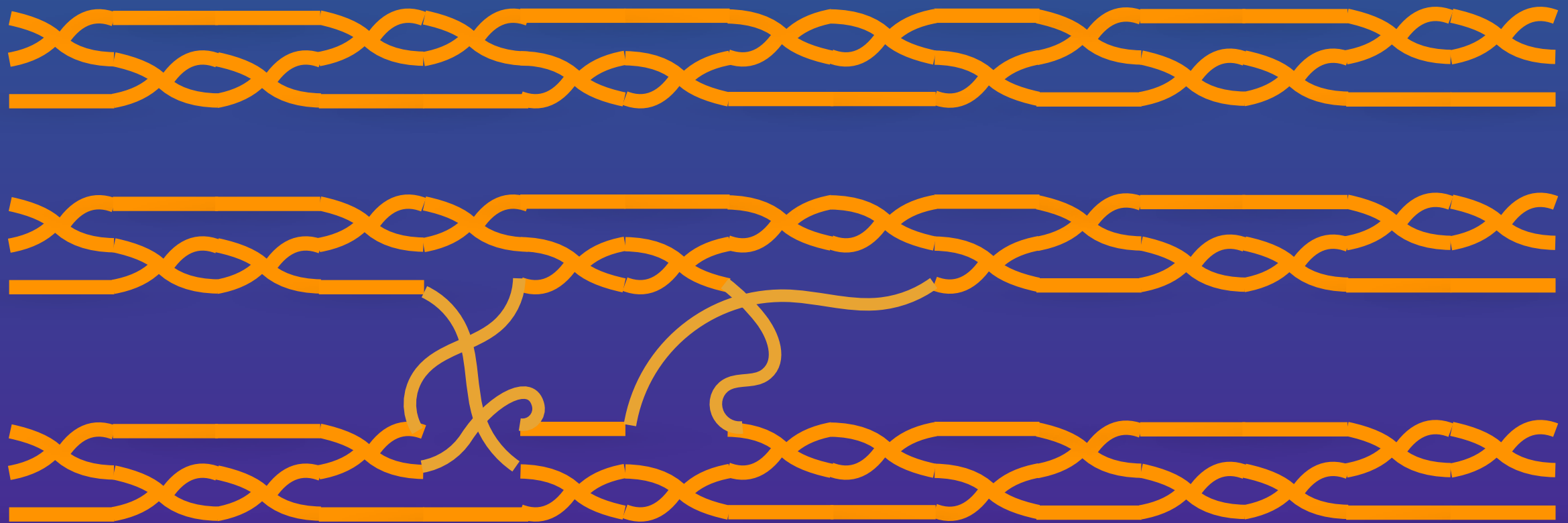
$\sigma_1 \sigma_2 \sigma_2 \sigma_1 \sigma_1$



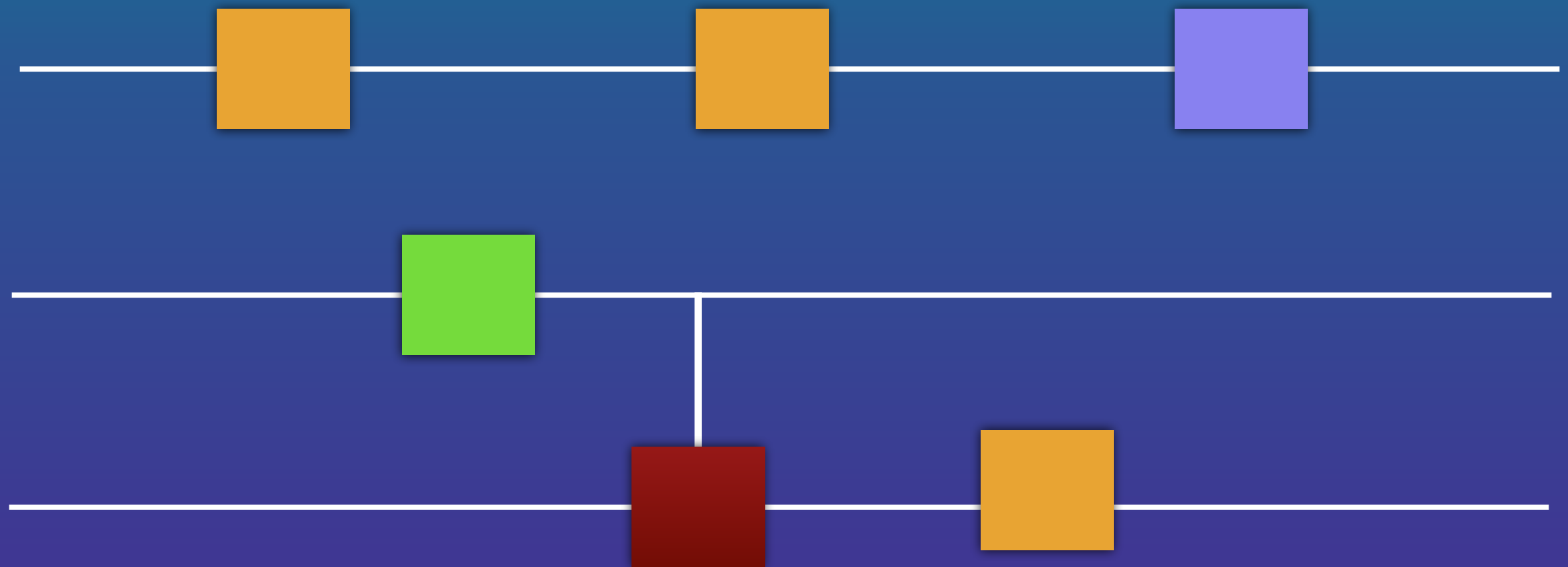
Two qubit gate Control Gate



With one qubit gate and a control two qubit gate,
we can do universal quantum computation.



With one qubit gate and a control two qubit gate,
we can do universal quantum computation.



[1] Braid Topologies for Quantum Computation

N. E. Bonesteel, Layla Hormozi, Georgios Zikos, Steven H. Simon, PhysRevLett.95.140503,

End of part II